

GLOBAL SOLUTIONS FOR 3D NONLOCAL GROSS-PITAEVSKII EQUATIONS WITH ROUGH DATA

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ABSTRACT. We study the Cauchy problem for the Gross-Pitaevskii equation with a nonlocal interaction potential of Hartree type in three space dimensions. If the potential is even and positive definite or a positive function and its Fourier transform decays sufficiently rapidly the problem is shown to be globally well-posed for large rough data which not necessarily have finite energy and also in a situation where the energy functional is not positive definite. The proof uses a suitable modification of the I-method.

1. INTRODUCTION AND MAIN RESULTS

We consider the Cauchy problem for the Gross-Pitaevskii equation with non-local nonlinearity in three space dimensions:

$$i \frac{\partial v}{\partial t} - \Delta v = v(W * (1 - |v|^2)) \quad (1)$$

$$v(x, 0) = v_0(x), \quad (2)$$

under the condition

$$v \rightarrow 1 \quad \text{as} \quad |x| \rightarrow +\infty, \quad (3)$$

where $v : \mathbb{R}^{1+3} \rightarrow \mathbb{C}$.

More generally one could also consider the condition

$$|v| \rightarrow 1 \quad \text{as} \quad |x| \rightarrow +\infty, \quad (4)$$

but for simplicity we restrict ourselves to (3). This problem was introduced by Gross [Gr] and Pitaevskii [P] in order to model the kinetic of a weakly interacting Bose gas. Here W describes the interaction between bosons. The original equation reads as follows

$$i \frac{\partial \psi}{\partial t}(x, t) + \frac{\hbar^2}{2m} \Delta \psi(x, t) = \psi(x, t) \int_{\mathbb{R}^n} V(x - y) |\psi(y, t)|^2 dy$$

and is equivalent modulo normalizations to equation (1), provided $W * 1$ is well-defined and positive, which in the cases we consider (under the assumptions (A1) and either (A2) or (A3) below) is obviously true. For a detailed derivation of our problem from the original Gross-Pitaevskii form we refer to [L].

The most studied case is $W = \delta$ (= Dirac function), which occurs in theoretical physics e.g. in superfluidity, nonlinear optics and Bose-Einstein condensation [C],[JPRo],[JR],[KL]. For further references we also refer to the introduction of [L].

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Using the energy conservation law which in the case $W = \delta$ is

$$E(v(t)) = \int (|\nabla v(x, t)|^2 + \frac{1}{2}(|v(x, t)|^2 - 1)^2) dx = E(v_0).$$

it was shown by Bethuel and Saut [BS], Appendix A, that the problem is globally well-posed for data of the form $v_0 \in 1 + H^1(\mathbb{R}^3)$. Gérard [Ge] proved the same result for data in the larger energy space in two and three space dimensions. Gallo [Ga] generalized these results to a class of local nonlinearities for data with finite energy and space dimension $n \leq 4$. The author [Pe] showed that global well-posedness holds true even for data with less regularity, namely $v_0 = 1 + u_0$, where $u_0 \in H^s(\mathbb{R}^3)$ for $5/6 < s < 1$. To prove this result one uses Bourgain type spaces and the so-called I -method (or method of almost conservation laws), which was introduced by Colliander, Keel, Staffilani, Takaoka and Tao [CKSTT] and successfully applied to various problems.

We now want to study the problem for two types of nonlocal nonlinearities. Nonlocal nonlinearities were as mentioned above already introduced by Gross and Pitaevskii. In the case of three space dimensions Shchesnovich and Kraenkel [SK] consider $W(x) = \frac{1}{4\pi\epsilon^2|x|} \exp(-\frac{|x|}{\epsilon})$ for $\epsilon > 0$ with Fourier transform $\widehat{W}(\xi) = \frac{1}{1+\epsilon^2|\xi|^2}$. The case $W = \chi_{\{|x| \leq a\}}$ (χ_A = characteristic function of the set A) was used in the study of supersolids [ABJ],[JPR],[PR]. These examples are included in the class of nonlocal nonlinearities with suitable mapping properties and positivity conditions on W considered by de Laire [L] such that the Cauchy problem (1),(2),(4) is globally well-posed in the space $\phi + H^1(\mathbb{R}^n)$, where ϕ has finite energy and fulfills suitable boundedness assumptions, in particular $|\phi(x)| \rightarrow 1$ as $|x| \rightarrow \infty$.

Our aim is to give similar results for less regular data. From now on we consider the case of three space dimensions and make the following

General Assumption on W :

(A1) $W : \mathbb{R}^3 \rightarrow \mathbb{R}$ is even, $W \in L^1(\mathbb{R}^3)$, $|\widehat{W}(\xi)| \lesssim \langle \xi \rangle^{-2} \forall \xi \in \mathbb{R}^3$

and **either**

(A2) $\widehat{W}(\xi) \geq 0 \quad \forall \xi \in \mathbb{R}^3$

or

(A3) $W(x) \geq 0 \quad \forall x \in \mathbb{R}^3$.

Let us remark, that \widehat{W} is real-valued and even, if W has the same properties.

We have especially the following two examples in mind, which we mentioned above:

Case A: $W(x) = \frac{1}{4\pi|x|} e^{-|x|}$.

We have $\widehat{W}(\xi) = \frac{1}{1+|\xi|^2}$, so that (A1),(A2) and (A3) are satisfied.

Case B: $W = \chi_{\{|x| \leq a\}}$.

Obviously (A3) is satisfied. We also have $\widehat{W}(\xi) = a^{-\frac{3}{2}} |\xi|^{-\frac{3}{2}} J_{\frac{3}{2}}(2\pi a |\xi|)$, where $J_{\frac{3}{2}}$ is the Bessel function of the first kind of order $\frac{3}{2}$, which has the properties $J_{\frac{3}{2}}(|\xi|) \sim |\xi|^{\frac{3}{2}}$ as $|\xi| \ll 1$ and $J_{\frac{3}{2}}(|\xi|) \lesssim \frac{1}{|\xi|^{\frac{1}{2}}}$ as $|\xi| \gg 1$. Thus (A1) is also satisfied.

For simplicity we assume $\phi \equiv 1$ and consider (1),(2),(3). As usual we transform the problem (1),(2),(3) by setting $u = v - 1$ into the equivalent form

$$i \frac{\partial u}{\partial t} - \Delta u + F(u) = 0 \tag{5}$$

$$u(x, 0) = u_0(x) \tag{6}$$

where

$$F(u) = (1 + u)(W * (|u|^2 + 2 \operatorname{Re} u)) \tag{7}$$

under the condition

$$u \rightarrow 0 \quad \text{as} \quad |x| \rightarrow +\infty. \quad (8)$$

Assuming W to be real-valued and even the conserved energy is given by

$$E(u(t)) = \int |\nabla u(t)|^2 dx + \frac{1}{2} \int (W * (|u|^2 + 2 \operatorname{Re} u)) (|u|^2 + 2 \operatorname{Re} u) dx. \quad (9)$$

Remark that no L^2 -conservation law holds.

Under our hypothesis on W de Laire's results [L] especially imply that the Cauchy problem (5),(6),(7),(8) is globally well-posed in $C^0(\mathbb{R}, H^1(\mathbb{R}^3))$ for data $u_0 \in H^1(\mathbb{R}^3)$. We now show that this problem for data $u_0 \in H^s(\mathbb{R}^3)$ is globally well-posed in $C^0(\mathbb{R}, H^s(\mathbb{R}^3))$, i.e. (1),(2),(3) for $v_0 \in 1 + H^s(\mathbb{R}^3)$, if $1/2 < s < 1$ by application of the I -method. As usual the energy conservation law is not directly applicable for H^s -data with $s < 1$. However there is an "almost conservation law" for the modified energy $E(Iu)$, which is well defined for $u \in H^s$ (see the definition of I below). If we assume (A1) and (A2), this leads to an a-priori bound of $\|\nabla Iu(t)\|_{L^2}$, if s is close enough to 1, namely $s > 1/2$, because the energy functional is positive definite, a property which is usually assumed when the I -method is applied. If we assume (A1) and (A3) however it is not obvious that the H^1 -norm of the solution can be controlled by the energy, because it is not definite. Nevertheless it is possible to modify the I -method in this case suitably, but the argument to get the required bound for $\|\nabla Iu(t)\|_{L^2}$ is more involved. Once a bound for $\|\nabla Iu(t)\|_{L^2}$ is achieved we can also deduce an a-priori bound for $\|u(t)\|_{L^2}$, which together gives an a-priori bound for $\|u(t)\|_{H^s}$.

The main results (cf. the definition of the $X^{s,b}$ -spaces below) are summarized in the following three theorems:

Theorem 1.1. (*Unconditional uniqueness*) Assume (A1) and moreover (A2) or (A3), $u_0 \in L^2(\mathbb{R}^3)$. The Cauchy problem (5),(6),(7) has at most one solution $u \in C^0([0, T], L^2(\mathbb{R}^3))$ for any $T > 0$.

Theorem 1.2. (*Local well-posedness*) Assume (A1) and moreover (A2) or (A3), $s \geq 0$ and $u_0 \in H^s(\mathbb{R}^3)$. Then the Cauchy problem (5),(6),(7) has a unique local solution $u \in X^{s, \frac{1}{2}+}[0, \delta]$, where δ can be chosen as $\delta \sim \|u_0\|_{H^s}^{-\frac{4}{2s+1}-}$. This solution belongs to $C^0([0, \delta], H^s(\mathbb{R}^3))$ and is also unique in this space.

Theorem 1.3. (*Global well-posedness*) Assume (A1) and moreover (A2) or (A3), $T > 0$, $s > 1/2$ and $u_0 \in H^s(\mathbb{R}^3)$. Then the Cauchy problem (5),(6),(7) has a unique global solution $u \in X^{s, \frac{1}{2}+}[0, T]$. This solution belongs to $C^0([0, T], H^s(\mathbb{R}^3))$ and is also unique in this space.

We use the following notation and well-known facts: the multiplier $I = I_N$ is for given $s < 1$ and $N \geq 1$ defined by

$$\widehat{I_N f}(\xi) := m_N(\xi) \widehat{f}(\xi),$$

where $\widehat{\cdot}$ denotes the Fourier transform with respect to the space variables. Here $m_N(\xi)$ is a smooth, radially symmetric, nonincreasing function of $|\xi|$ with

$$m_N(\xi) = \begin{cases} 1 & |\xi| \leq N \\ (\frac{N}{|\xi|})^{1-s} & |\xi| \geq 2N \end{cases}$$

We remark that $I : H^s \rightarrow H^1$ is a smoothing operator, so that especially $E(Iu)$ is well-defined for $u \in H^s(\mathbb{R}^3)$. This follows from $W \in L^1$, Young's inequality and Sobolev's embedding $H^1(\mathbb{R}^3) \subset L^4(\mathbb{R}^3)$.

We use the Bourgain type function space $X^{m,b}$ belonging to the Schrödinger equation $iu_t - \Delta u = 0$, which is defined as follows: let $\widehat{\cdot}$ or \mathcal{F} denote the Fourier

transform with respect to space and time and \mathcal{F}^{-1} its inverse. $X^{m,b}$ is the completion of $\mathcal{S}(\mathbb{R} \times \mathbb{R}^3)$ with respect to

$$\|f\|_{X^{m,b}} = \|\langle \xi \rangle^m \langle \tau \rangle^b \mathcal{F}(e^{-it\Delta} f(x, t))\|_{L_{\xi, \tau}^2} = \|\langle \xi \rangle^m \langle \tau + |\xi|^2 \rangle^b \widehat{f}(\xi, \tau)\|_{L_{\xi, \tau}^2},$$

For a given time interval I we define

$$\|f\|_{X^{m,b}(I)} := \inf_{g|_I = f} \|g\|_{X^{m,b}}.$$

We recall the following facts about the solutions u of the inhomogeneous linear Schrödinger equation (see e.g. [GTV])

$$iu_t - \Delta u = F, \quad u(0) = f. \quad (10)$$

For $b' + 1 \geq b \geq 0 \geq b' > -1/2$ and $T \leq 1$ we have

$$\|u\|_{X^{m,b}[0,T]} \lesssim \|f\|_{H^m} + T^{1+b'-b} \|F\|_{X^{m,b'}[0,T]}.$$

For $1/2 > b > b' \geq 0$ or $0 \geq b > b' > -1/2$:

$$\|f\|_{X^{m,b'}[0,T]} \lesssim T^{b-b'} \|f\|_{X^{m,b}[0,T]}$$

(see e.g. [G], Lemma 1.10).

Fundamental are the following Strichartz type estimates for the solution u of (10) in three space dimensions (see [CH],[KT]):

$$\|u\|_{L^q(I, L^r(\mathbb{R}^3))} \lesssim \|f\|_{L^2(\mathbb{R}^3)} + \|F\|_{L^{\tilde{q}'}(I, L^{\tilde{r}'}(\mathbb{R}^3))}$$

with implicit constant independent of the interval $I \subset \mathbb{R}$ for all pairs $(q, r), (\tilde{q}, \tilde{r})$ with $q, r, \tilde{q}, \tilde{r} \geq 2$ and $\frac{1}{q} + \frac{3}{2r} = \frac{3}{4}$, $\frac{1}{\tilde{q}} + \frac{3}{2\tilde{r}} = \frac{3}{4}$, where $\frac{1}{q} + \frac{1}{q'} = 1$ and $\frac{1}{\tilde{r}} + \frac{1}{\tilde{r}'} = 1$. This implies

$$\|\psi\|_{L^q(I, L^r(\mathbb{R}^3))} \lesssim \|\psi\|_{X^{0, \frac{1}{2}+}(I)}.$$

For real numbers a we denote by $a+$, $a++$, $a-$ and $a--$ the numbers $a+\epsilon$, $a+2\epsilon$, $a-\epsilon$ and $a-2\epsilon$, respectively, where $\epsilon > 0$ is sufficiently small. We also use the notation $\langle x \rangle := (1 + |x|^2)^{\frac{1}{2}}$ for $x \in \mathbb{R}^3$.

The paper is organized as follows: in section 1 we prove the uniqueness result Theorem 1.1 and two versions of a local well-posedness result for (5),(6),(7), namely $u \in X^{s, \frac{1}{2}+}[0, \delta]$ for data $u_0 \in H^s$ with $s \geq 0$ (Theorem 1.2), and a modification where $\nabla Iu \in X^{0, \frac{1}{2}+}[0, \delta]$ for data $\nabla Iu_0 \in L^2$ (Proposition 2.1), which is necessary in order to combine it with an almost conservation law for the modified energy $E(Iu)$. In section 2 we use these local results and bounds for the modified energy given in section 3 in order to get the main theorem (Theorem 1.3). Under the assumptions (A1) and (A2) it is namely shown that the bounds for the modified energy immediately give a polynomial bound for $\|\nabla Iu(t)\|_{L^2}$, which can be shown to imply a uniform exponential bound for $\|u(t)\|_{L^2}$, and as a consequence for $\|u(t)\|_{H^s}$, which in view of the local well-posedness result suffices to get a global solution. Under the assumptions (A1) and (A3) we cannot immediately get a bound for $\|\nabla Iu(t)\|_{L^2}$ from the bound for the modified energy, but first we have to show an (exponential) bound for $\|Iu(t)\|_{L^2}$, which together with an energy bound gives the desired (exponential) bound for $\|\nabla Iu(t)\|_{L^2}$ and after that as in the previous case the global well-posedness result. In section 3 we first calculate $\frac{d}{dt}E(Iu)$ for any solution of the equation (5) and estimate the time integrated terms which appear in $\frac{d}{dt}E(Iu)$, which is the most complicated part. In section 2 these estimates are shown to control the modified energy $E(Iu)$ uniformly on arbitrary time intervals $[0, T]$, provided $s > 1/2$

2. UNIQUENESS AND LOCAL WELL-POSEDNESS

Proof of Theorem 1.1. Let $u, v \in C^0([0, T], L^2)$ be two solutions. Using Strichartz type estimates in order to control

$$\|u - v\|_{L_t^2 L_x^6} + \|u - v\|_{L_t^\infty L_x^2}$$

we have to estimate the various terms of $F(u) - F(v)$. By (A1) and the Hausdorff-Young inequality we have $W \in L^p$ for $1 \leq p < 3$ so that by Young's inequality we get

$$\begin{aligned} \|(W * |u|^2)(u - v)\|_{L_t^{1+} L_x^{2-}} &\lesssim \|W * |u|^2\|_{L_t^{2+} L_x^{3-}} \|u - v\|_{L_t^2 L_x^6} \\ &\lesssim \|W\|_{L_x^{3-}} \|u\|_{L_t^{4+} L_x^2}^2 \|u - v\|_{L_t^2 L_x^6} \\ &\lesssim T^{\frac{1}{2}-} \|u\|_{L_t^\infty L_x^2}^2 \|u - v\|_{L_t^2 L_x^6} \\ \|(W * (|u|^2 - |v|^2))v\|_{L_t^{1+} L_x^{2-}} &\lesssim \|W * (|u|^2 - |v|^2)\|_{L_t^2 L_x^\infty} \|u\|_{L_t^{2+} L_x^2} \\ &\lesssim \|W\|_{L_x^{3-}} \| |u|^2 - |v|^2 \|_{L_t^2 L_x^{\frac{3}{2}}} \|u\|_{L_t^{2+} L_x^2} \\ &\lesssim T^{\frac{1}{2}-} (\|u\|_{L_t^\infty L_x^2} + \|v\|_{L_t^\infty L_x^2}) \|u - v\|_{L_t^2 L_x^6} \|u\|_{L_t^\infty L_x^2} \\ \|(W * Re u)(u - v)\|_{L_t^1 L_x^2} &\lesssim \|W * Re u\|_{L_t^1 L_x^\infty} \|u - v\|_{L_t^\infty L_x^2} \\ &\lesssim \|W\|_{L^2 T} \|u\|_{L_t^\infty L_x^2} \|u - v\|_{L_t^\infty L_x^2} \\ \|W * Re(u - v)\|_{L_t^1 L_x^2} &\lesssim \|W\|_{L^1 T} \|u\|_{L_t^\infty L_x^2} \|u - v\|_{L_t^\infty L_x^2} \end{aligned}$$

Similarly the remaining terms can be estimated. Therefore, choosing T small enough, we get $u \equiv v$. \square

Next we prove the local well-posedness results.

Proof of Theorem 1.2. We want to apply the contraction mapping principle in the Bourgain type space $X^{s, \frac{1}{2}+}[0, \delta]$. We have to estimate

$$\|(W * (|u|^2 + 2Re u))(1 + u)\|_{X^{s, -\frac{1}{2}++}},$$

We denote by D^l the operator with symbol $|\xi|^l$ and similarly by $\langle D \rangle^l$ the operator with symbol $\langle \xi \rangle^l$. In order to estimate the cubic term by $\delta^{\frac{1}{2}+s-} \|u\|_{X^{s, \frac{1}{2}+}}^3$ we have to show (ignoring complex conjugates, which play no role for the calculations):

$$\int \langle D \rangle^s (\langle D \rangle^{-2} (u_1 u_2) u_3) \psi dx dt \lesssim \delta^{\frac{1}{2}+s-} \prod_{i=1}^3 \|u_i\|_{X^{s, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}.$$

This is sufficient, because we can assume w.l.o.g. that the Fourier transforms $\widehat{u}_i(\xi_i, \tau_i)$ and $\widehat{\psi}(\xi, \tau)$ are nonnegative, so that using the fundamental assumption $|\widehat{W}(\xi)| \lesssim \langle \xi \rangle^{-2}$ it is possible to replace here and in similar situations in the following the convolution with W by application of $\langle D \rangle^{-2}$. Using the Leibniz rule for fractional derivatives we reduce to the estimates (assuming w.l.o.g. $|\xi_2| \geq |\xi_1|$):

$$\|\langle D \rangle^{-2} (u_1 \langle D \rangle^s u_2) u_3 \psi\|_{L_{xt}^1} \lesssim \delta^{\frac{1}{2}+s-} \prod_{i=1}^3 \|u_i\|_{X^{s, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}$$

and

$$\|\langle D \rangle^{-2} (u_1 u_2) \langle D \rangle^s u_3 \psi\|_{L_{xt}^1} \lesssim \delta^{\frac{1}{2}+s-} \prod_{i=1}^3 \|u_i\|_{X^{s, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}.$$

We get

$$\|\langle D \rangle^{-2} (u_1 \langle D \rangle^s u_2) u_3 \psi\|_{L_{xt}^1} \lesssim \|\langle D \rangle^{-2} (u_1 \langle D \rangle^s u_2) u_3\|_{L_t^{1+} L_x^2} \|\psi\|_{L_t^\infty L_x^2}$$

and

$$\begin{aligned}
\|\langle D \rangle^{-2} (u_1 \langle D \rangle^s u_2) u_3\|_{L_t^{1+} L_x^2} &\lesssim \|\langle D \rangle^{-2} (u_1 \langle D \rangle^s u_2)\|_{L_t^{1+} L_x^{\hat{p}}} \|u_3\|_{L_t^\infty L_x^{\hat{q}}} \\
&\lesssim \|u_1 \langle D \rangle^s u_2\|_{L_t^{1+} L_x^{\hat{p}}} \|u_3\|_{L_t^\infty H_x^s} \\
&\lesssim \|u_1\|_{L_t^{1+} L_x^{\hat{p}}} \|\langle D \rangle^s u_2\|_{L_t^\infty L_x^2} \|u_3\|_{L_t^\infty H_x^s} \\
&\lesssim \|u_1\|_{L_t^{1+} H_x^{s, \hat{t}}} \|u_2\|_{X^{s, \frac{1}{2}+}} \|u_3\|_{X^{s, \frac{1}{2}+}} \\
&\lesssim \delta^{s+\frac{1}{2}-} \|u_1\|_{L_t^{\hat{w}} H_x^{s, \hat{t}}} \|u_2\|_{X^{s, \frac{1}{2}+}} \|u_3\|_{X^{s, \frac{1}{2}+}} \\
&\lesssim \delta^{s+\frac{1}{2}-} \prod_{i=1}^3 \|u_i\|_{X^{s, \frac{1}{2}+}}.
\end{aligned}$$

Here we set $\frac{1}{\hat{q}} = \frac{1}{2} - \frac{s}{3}$, $\frac{1}{\hat{p}} = \frac{s}{3}$, $\frac{1}{\hat{r}} = \frac{2}{3} + \frac{s}{3}$, so that $H_x^s \subset L_x^{\hat{q}}$, and $\frac{1}{\hat{v}} = \frac{1}{6} + \frac{s}{3}$, $\frac{1}{\hat{t}} = \frac{1}{6} + \frac{2}{3}s$, so that $H_x^{s, \hat{t}} \subset L_x^{\hat{v}}$, and finally $\frac{1}{\hat{w}} = \frac{1}{2} - s$, so that $\frac{1}{\hat{w}} + \frac{3}{2\hat{t}} = \frac{3}{4}$, which allows to apply Strichartz' estimate in the last line. Similarly we get

$$\begin{aligned}
\|\langle D \rangle^{-2} (u_1 u_2) \langle D \rangle^s u_3\|_{L_t^{1+} L_x^2} &\lesssim \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^{1+} L_x^\infty} \|\langle D \rangle^s u_3\|_{L_t^\infty L_x^2} \\
&\lesssim \|u_1 u_2\|_{L_t^{1+} L_x^{\frac{3}{2}+}} \|u_3\|_{X^{s, \frac{1}{2}+}} \\
&\lesssim \|u_1\|_{L_t^{1+} L_x^{\hat{v}}} \|u_2\|_{L_t^\infty L_x^{\hat{q}}} \|u_3\|_{X^{s, \frac{1}{2}+}} \\
&\lesssim \delta^{s+\frac{1}{2}-} \|u_1\|_{L_t^{\hat{w}} H_x^{s, \hat{t}}} \|u_2\|_{L_t^\infty H_x^s} \|u_3\|_{X^{s, \frac{1}{2}+}} \\
&\lesssim \delta^{s+\frac{1}{2}-} \prod_{i=1}^3 \|u_i\|_{X^{s, \frac{1}{2}+}}.
\end{aligned}$$

Here $\frac{1}{\hat{q}} = \frac{1}{2} - \frac{s}{3}$, $\frac{1}{\hat{v}} = \frac{1}{6} + \frac{s}{3}$, $\frac{1}{\hat{t}} = \frac{1}{6} + \frac{2}{3}s$, $\frac{1}{\hat{w}} = \frac{1}{2} - s$, so that $\frac{1}{\hat{w}} + \frac{3}{2\hat{t}} = \frac{3}{4}$ allows to apply Strichartz' estimate again.

The quadratic terms are handled as follows (assuming again $|\xi_2| \geq |\xi_1|$):

$$\begin{aligned}
\|\langle D \rangle^{-2+s} (u_1 u_2)\|_{L_t^{1+} L_x^2} &\lesssim \|\langle D \rangle^s (u_1 u_2)\|_{L_t^{1+} L_x^1} \lesssim \|u_1 \langle D \rangle^s u_2\|_{L_t^{1+} L_x^1} \\
&\lesssim \|u_1\|_{L_t^{2+} L_x^2} \|\langle D \rangle^s u_2\|_{L_t^2 L_x^2} \lesssim \delta^{1-} \|u_1\|_{X^{s, \frac{1}{2}+}} \|u_2\|_{X^{s, \frac{1}{2}+}}
\end{aligned}$$

and

$$\begin{aligned}
&\|\langle D \rangle^s ((\langle D \rangle^{-2} u_1) u_2)\|_{L_t^{1+} L_x^2} \\
&\lesssim \|(\langle D \rangle^{-2+s} u_1) u_2\|_{L_t^{1+} L_x^2} + \|\langle D \rangle^{-2} u_1 \langle D \rangle^s u_2\|_{L_t^{1+} L_x^2} \\
&\lesssim \|\langle D \rangle^{-2+s} u_1\|_{L_t^{1+} L_x^\infty} \|u_2\|_{L_t^\infty L_x^2} + \|\langle D \rangle^{-2} u_1\|_{L_{t,x}^\infty} \|\langle D \rangle^s u_2\|_{L_t^{1+} L_x^2} \\
&\lesssim \delta^{1-} \|\langle D \rangle^s u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2} + \delta^{1-} \|u_1\|_{L_t^\infty L_x^2} \|\langle D \rangle^s u_2\|_{L_t^\infty L_x^2} \\
&\lesssim \delta^{1-} \|u_1\|_{X^{s, \frac{1}{2}+}} \|u_2\|_{X^{s, \frac{1}{2}+}}.
\end{aligned}$$

Similar estimates hold for the difference $F(u) - F(v)$, so that a standard Picard iteration under the conditions $\delta^{s+\frac{1}{2}-} \|u_0\|_{H^s}^2 \lesssim 1$ and $\delta^{1-} \|u_0\|_{H^s} \lesssim 1$ shows the existence of a unique solution in $X^{s, \frac{1}{2}+}[0, \delta] \subset C^0([0, \delta], H^s)$. It is also unique in this latter space by Theorem 1.1. \square

Remark: This Theorem shows that in order to get a global solution it is sufficient to give an a-priori bound of $\|u(t)\|_{H^s}$.

We next prove a modified local well-posedness result involving the operator I (recall that I depends on s and N).

Proposition 2.1. *Assume $s \geq 0$ and $\nabla I u_0 \in L^2$. Then the Cauchy problem (5),(6),(7) (after application of I) has a unique local solution u with $\nabla I u \in$*

$X^{0, \frac{1}{2}+}[0, \delta]$ and $\|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]} \leq \sqrt{M}\|\nabla Iu_0\|_{L^2}$, where $M \geq 1$ is independent of u_0 , and $\delta \leq 1$ can be chosen such that

$$(\delta^{\frac{1}{2}-}N^{-2} + \delta^{1-})\|\nabla Iu_0\|_{L^2}^2 \sim 1.$$

Proof. The cubic term in the nonlinearity will be estimated as follows (dropping $[0, \delta]$ from the notation):

$$\|\nabla I(W * (u_1 u_2) u_3)\|_{X^{0, -\frac{1}{2}++}} \lesssim (\frac{\delta^{\frac{1}{2}-}}{N^2} + \delta^{1-}) \prod_{i=1}^3 \|\nabla Iu_i\|_{X^{0, \frac{1}{2}+}}.$$

This follows from

$$\begin{aligned} A &:= \int_0^\delta \int_* M(\xi_1, \xi_2, \xi_3) \prod_{i=1}^3 \widehat{u}_i(\xi_i, t) \widehat{\psi}(\xi_4, t) d\xi_1 d\xi_2 d\xi_3 d\xi_4 dt \\ &\lesssim (\frac{\delta^{\frac{1}{2}-}}{N^2} + \delta^{1-}) \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}+}}, \end{aligned}$$

where

$$M(\xi_1, \xi_2, \xi_3) := \frac{m(\xi_1 + \xi_2 + \xi_3)}{m(\xi_1)m(\xi_2)m(\xi_3)} \cdot \frac{|\xi_1 + \xi_2 + \xi_3|}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1| |\xi_2| |\xi_3|},$$

and $*$ denotes integration over the region $\{\sum_{i=1}^4 \xi_i = 0\}$. We assume here and in the following again w.l.o.g. that the Fourier transforms are nonnegative, and also w.l.o.g. that $|\xi_1| \geq |\xi_2|$. We again used the property $|\widehat{W}(\xi)| \lesssim \langle \xi \rangle^{-2}$.

We make a case by case analysis depending on the relative size of the frequencies.

Case 1: $|\xi_1| \geq |\xi_2| \geq |\xi_3| \gtrsim N$.

1a. $|\xi_1 + \xi_2 + \xi_3| \geq N$.

We get

$$\begin{aligned} M(\xi_1, \xi_2, \xi_3) &\lesssim \prod_{i=1}^3 (\frac{|\xi_i|}{N})^{1-s} \frac{N^{1-s}}{|\xi_1 + \xi_2 + \xi_3|^{1-s}} \frac{|\xi_1 + \xi_2 + \xi_3|}{|\xi_1| |\xi_2| |\xi_3| \langle \xi_1 + \xi_2 \rangle^2} \\ &\lesssim \frac{1}{N^{2(1-s)}} \frac{|\xi_1|^s}{|\xi_1|^s |\xi_2|^s |\xi_3|^s \langle \xi_1 + \xi_2 \rangle^2} \lesssim \frac{1}{N^2 \langle \xi_1 + \xi_2 \rangle^2}. \end{aligned}$$

This implies by Sobolev's embedding and Strichartz' estimates:

$$\begin{aligned} A &\lesssim \frac{1}{N^2} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^3} \|u_3\|_{L_t^2 L_x^6} \|\psi\|_{L_t^2 L_x^2} \\ &\lesssim \frac{1}{N^2} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^2 L_x^6} \delta^{\frac{1}{2}-} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{1}{N^2} \delta^{\frac{1}{2}-} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}. \end{aligned}$$

1b. $|\xi_1 + \xi_2 + \xi_3| \leq N$.

We have

$$\begin{aligned} M(\xi_1, \xi_2, \xi_3) &\lesssim \prod_{i=1}^3 (\frac{|\xi_i|}{N})^{1-s} \frac{N}{|\xi_1| |\xi_2| |\xi_3| \langle \xi_1 + \xi_2 \rangle^2} \\ &\lesssim \frac{N}{N^{3(1-s)} |\xi_1|^s |\xi_2|^s |\xi_3|^s \langle \xi_1 + \xi_2 \rangle^2} \lesssim \frac{1}{N^2 \langle \xi_1 + \xi_2 \rangle^2} \end{aligned}$$

as in case 1a.

Case 2: $|\xi_3| \geq |\xi_1| \geq |\xi_2|$.

This case can be treated similarly as case 1.

Case 3: $|\xi_1| \geq |\xi_2| \gtrsim N \geq |\xi_3|$.

3a. $|\xi_1 + \xi_2 + \xi_3| \geq N$.

We get

$$\begin{aligned} M(\xi_1, \xi_2, \xi_3) &\lesssim \frac{|\xi_1|^{1-s}}{N^{1-s}} \frac{|\xi_2|^{1-s}}{N^{1-s}} \frac{N^{1-s}}{|\xi_1 + \xi_2 + \xi_3|^{1-s}} \frac{|\xi_1 + \xi_2 + \xi_3|}{|\xi_1||\xi_2||\xi_3|(\xi_1 + \xi_2)^2} \\ &\lesssim \frac{|\xi_1|^s}{|\xi_1|^s |\xi_2|^s |\xi_3| N^{1-s} \langle \xi + \xi_2 \rangle^2} \lesssim \frac{1}{N |\xi_3| \langle \xi_1 + \xi_2 \rangle^2}, \end{aligned}$$

so that

$$\begin{aligned} A &\lesssim \frac{1}{N} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^3} \|D^{-1} u_3\|_{L_t^\infty L_x^6} \|\psi\|_{L_t^1 L_x^2} \\ &\lesssim \frac{1}{N} \delta^{1-} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{1}{N} \delta^{1-} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}. \end{aligned}$$

3b. $|\xi_1 + \xi_2 + \xi_3| \leq N$.

Similarly we get

$$M(\xi_1, \xi_2, \xi_3) \lesssim \frac{|\xi_1|^{1-s}}{N^{1-s}} \frac{|\xi_2|^{1-s}}{N^{1-s}} \frac{N}{|\xi_1||\xi_2||\xi_3|(\xi_1 + \xi_2)^2} \lesssim \frac{1}{N |\xi_3| \langle \xi_1 + \xi_2 \rangle^2}$$

as in case 3a.

Case 4: $|\xi_1|, |\xi_3| \gtrsim N \gtrsim |\xi_2|$.

Similarly as in case 3 we get

$$M(\xi_1, \xi_2, \xi_3) \lesssim \frac{1}{N |\xi_2| \langle \xi_1 + \xi_2 \rangle^2},$$

so that by Sobolev's embedding and Strichartz' estimates

$$\begin{aligned} A &\lesssim \frac{1}{N} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^\infty} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^1 L_x^2} \\ &\lesssim \frac{1}{N} \delta^{1-} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{1}{N} \delta^{1-} \|u_1\|_{L_t^\infty L_x^{2+}} \|D^{-1} u_2\|_{L_t^\infty L_x^6} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{1}{N} \delta^{1-} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}. \end{aligned}$$

Case 5: $|\xi_3| \gtrsim N \gg |\xi_1| \geq |\xi_2|$.

We have

$$\begin{aligned} M(\xi_1, \xi_2, \xi_3) &\lesssim \frac{|\xi_3|^{1-s}}{N^{1-s}} \frac{N^{1-s}}{|\xi_1 + \xi_2 + \xi_3|^{1-s}} \frac{|\xi_1 + \xi_2 + \xi_3|}{|\xi_1||\xi_2||\xi_3|(\xi_1 + \xi_2)^2} \\ &\lesssim \frac{1}{|\xi_1||\xi_2| \langle \xi_1 + \xi_2 \rangle^2}, \end{aligned}$$

leading to

$$\begin{aligned} A &\lesssim \|\langle D \rangle^{-2} (D^{-1} u_1 D^{-1} u_2)\|_{L_t^\infty L_x^\infty} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^1 L_x^2} \\ &\lesssim \delta^{1-} \|D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^3} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \delta^{1-} \|D^{-1} u_1\|_{L_t^\infty L_x^6} \|D^{-1} u_2\|_{L_t^\infty L_x^6} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \delta^{1-} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}. \end{aligned}$$

Case 6: $|\xi_1| \gtrsim N \gg |\xi_2|, |\xi_3|$.
Similarly as in case 5 we get

$$M(\xi_1, \xi_2, \xi_3) \lesssim \frac{1}{|\xi_2||\xi_3|\langle \xi_1 + \xi_2 \rangle^2},$$

which implies

$$\begin{aligned} A &\lesssim \delta^{1-} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^3} \|D^{-1} u_3\|_{L_t^\infty L_x^6} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \delta^{1-} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}}} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \delta^{1-} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^6} \|u_3\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\ &\lesssim \delta^{1-} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}. \end{aligned}$$

Case 7: $N \gg |\xi_1|, |\xi_2|, |\xi_3|$.
We easily get

$$M(\xi_1, \xi_2, \xi_3) \lesssim \frac{|\xi_1 + \xi_2 + \xi_3|}{|\xi_1||\xi_2||\xi_3|\langle \xi_1 + \xi_2 \rangle^2} \lesssim \frac{1}{|\xi_2||\xi_3|\langle \xi_1 + \xi_2 \rangle^2} \text{ or } \lesssim \frac{1}{|\xi_1||\xi_2|\langle \xi_1 + \xi_2 \rangle^2},$$

which can be handled like case 6 or case 5. This completes the claimed estimate for the cubic term.

Next we consider the quadratic terms in the nonlinearity. They turn out to be less critical. First we prove the estimate

$$\|\nabla IW * (u_1 u_2)\|_{X^{0, -\frac{1}{2}++}} \lesssim \left(\frac{\delta^{\frac{1}{2}-}}{N} + \delta^{1-}\right) \|\nabla I u_1\|_{X^{0, \frac{1}{2}+}} \|\nabla I u_2\|_{X^{0, \frac{1}{2}+}}.$$

This follows from

$$\begin{aligned} B &:= \int_0^\delta \int_* M(\xi_1, \xi_2) \prod_{i=1}^2 \widehat{u}_i(\xi_i, t) \widehat{\psi}(\xi_3, t) d\xi_1 d\xi_2 d\xi_3 dt \\ &\lesssim \left(\frac{\delta^{\frac{1}{2}-}}{N} + \delta^{1-}\right) \prod_{i=1}^2 \|u_i\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}+}}, \end{aligned}$$

where

$$M(\xi_1, \xi_2) := \frac{m(\xi_1 + \xi_2)}{m(\xi_1)m(\xi_2)} \cdot \frac{|\xi_1 + \xi_2|}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1| |\xi_2|},$$

and $*$ denotes integration over the region $\{\sum_{i=1}^3 \xi_i = 0\}$. We assume w.l.o.g. $|\xi_1| \geq |\xi_2|$.

Case 1: $|\xi_1| \geq |\xi_2| \geq N$.

1a. $|\xi_1 + \xi_2| \geq N$.

We get

$$\begin{aligned} M(\xi_1, \xi_2) &\lesssim \frac{|\xi_1|^{1-s}}{N^{1-s}} \frac{|\xi_2|^{1-s}}{N^{1-s}} \frac{N^{1-s}}{|\xi_1 + \xi_2|^{1-s}} \frac{|\xi_1 + \xi_2|}{|\xi_1||\xi_2|\langle \xi_1 + \xi_2 \rangle^2} \\ &\lesssim \frac{|\xi_1 + \xi_2|^s}{|\xi_1|^s |\xi_2|^s N^{1-s} \langle \xi_1 + \xi_2 \rangle^2} \lesssim \frac{1}{N \langle \xi_1 + \xi_2 \rangle^2}, \end{aligned}$$

so that by Strichartz' estimates

$$\begin{aligned}
B &\lesssim \frac{1}{N} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^2 L_x^2} \|\psi\|_{L_t^2 L_x^2} \\
&\lesssim \frac{1}{N} \|u_1 u_2\|_{L_t^2 L_x^{\frac{3}{2}}} \|\psi\|_{L_t^2 L_x^2} \\
&\lesssim \frac{1}{N} \delta^{\frac{1}{2}-} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^2 L_x^6} \|\psi\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{1}{N} \delta^{\frac{1}{2}-} \|u_1\|_{X^{0, \frac{1}{2}+}} \|u_2\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}.
\end{aligned}$$

1b. $|\xi_1 + \xi_2| \leq N$.

This case is similar to case 1a.

Case 2: $N \gtrsim |\xi_2|$ and $|\xi_1| \gg |\xi_2|$.

One has

$$M(\xi_1, \xi_2) \lesssim \frac{|\xi_1 + \xi_2|}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1| |\xi_2|} \lesssim \frac{1}{|\xi_2| \langle \xi_1 + \xi_2 \rangle^2},$$

so that

$$\begin{aligned}
B &\lesssim \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^2 L_x^2} \|\psi\|_{L_t^2 L_x^2} \\
&\lesssim \|u_1 D^{-1} u_2\|_{L_t^2 L_x^{\frac{3}{2}}} \|\psi\|_{L_t^2 L_x^2} \\
&\lesssim \delta^{1-} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^6} \|\psi\|_{L_t^\infty L_x^2} \\
&\lesssim \delta^{1-} \|u_1\|_{X^{0, \frac{1}{2}+}} \|u_2\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}.
\end{aligned}$$

Case 3: $N \geq |\xi_1| \sim |\xi_2|$.

This case can be handled like case 2.

Finally we show

$$\|\nabla I((W * u_1) u_2)\|_{X^{0, -\frac{1}{2}++}} \lesssim \delta^{1-} \|\nabla I u_1\|_{X^{0, \frac{1}{2}+}} \|\nabla I u_2\|_{X^{0, \frac{1}{2}+}}.$$

Here B is as in case 2 with

$$M(\xi_1, \xi_2) := \frac{m(\xi_1 + \xi_2)}{m(\xi_1) m(\xi_2)} \cdot \frac{|\xi_1 + \xi_2|}{\langle \xi_1 \rangle^2 |\xi_1| |\xi_2|}.$$

Because the estimates are similar to the previous case we only consider the most critical low frequency cases.

Case 1: $|\xi_1| \gg |\xi_2|$ and $N \geq |\xi_1|$ (or $N \gg |\xi_1| \sim |\xi_2|$).

The estimate

$$M(\xi_1, \xi_2) \lesssim \frac{1}{\langle \xi_1 \rangle^2 |\xi_1|}$$

implies

$$\begin{aligned}
B &\lesssim \|\langle D \rangle^{-2} D^{-1} u_1\|_{L_{x,t}^\infty} \|u_2\|_{L_t^2 L_x^2} \|\psi\|_{L_t^2 L_x^2} \\
&\lesssim \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^2 L_x^2} \|\psi\|_{L_t^2 L_x^2} \\
&\lesssim \delta^{1-} \|u_1\|_{X^{0, \frac{1}{2}+}} \|u_2\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}.
\end{aligned}$$

Case 2: $|\xi_1| \gg |\xi_2|$ and $N \geq |\xi_2|$.

The estimate

$$M(\xi_1, \xi_2) \lesssim \frac{1}{\langle \xi_1 \rangle^2 |\xi_2|}$$

implies

$$\begin{aligned}
B &\lesssim \|\langle D \rangle^{-2} u_1\|_{L_t^\infty L_x^3} \|D^{-1} u_2\|_{L_t^2 L_x^6} \|\psi\|_{L_t^2 L_x^2} \\
&\lesssim \delta^{1-} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2} \|\psi\|_{L_t^\infty L_x^2} \\
&\lesssim \delta^{1-} \|u_1\|_{X^{0, \frac{1}{2}+}} \|u_2\|_{X^{0, \frac{1}{2}+}} \|\psi\|_{X^{0, \frac{1}{2}-}}.
\end{aligned}$$

We remark that similar estimates can be given for the difference terms in order to use Banach's fixed point theorem. In order to get a contraction we have to fulfill the estimates

$$(\delta^{\frac{1}{2}-} N^{-2} + \delta^{1-}) \|\nabla I u_0\|_{L^2}^2 \ll 1 \quad \text{and} \quad (\delta^{\frac{1}{2}-} N^{-1} + \delta^{1-}) \|\nabla I u_0\|_{L^2} \ll 1.$$

The latter requirement is weaker, so that the claimed result follows. \square

Remark: We want to iterate this local existence theorem with time steps of equal length until we reach a given (large) time T . To achieve this we need to control

$$\|\nabla I u(t)\|_{L^2} \leq c(T) \quad \forall 0 \leq t \leq T. \quad (11)$$

This will be shown under the assumption $u_0 \in H^s$ with $s > 1/2$.

3. PROOF OF THEOREM 1.3

In this section we first show that the bound (11) implies global well-posedness and after that we derive such a bound from the estimates for the modified energy $E(Iu)$ in the next section.

Proof of Theorem 1.3. So let us assume for the moment that (11) holds. This means that on any existence interval $[0, T]$ we have an a-priori bound (for fixed N) of

$$\|\nabla I u(t)\|_{L^2} \sim \| |\xi| \widehat{u}(\xi, t) \|_{L^2(\{|\xi| \leq N\})} + \| |\xi|^s \widehat{u}(\xi, t) \|_{L^2(\{|\xi| \geq N\})} N^{1-s}, \quad (12)$$

especially

$$\| |\xi|^s \widehat{u}(\xi, t) \|_{L^2(\{|\xi| \geq 1\})} \leq c(T). \quad (13)$$

If we can show that this implies an a-priori bound for $\|u(t)\|_{L^2}$, which is done in the following lemma, we immediately get an a-priori bound for $\|u(t)\|_{H^s}$, $0 \leq t \leq T$, thus a unique global solution in $X^{s, \frac{1}{2}+}[0, T] \subset C^0([0, T], H^s)$ for any T using our local well-posedness result (Theorem 1.2), which is also unique in this latter space by Theorem 1.1. \square

Lemma 3.1. *Assume (11) and $s \geq 1/2$. On any existence interval $[0, T]$ of the solution $u \in X^{s, \frac{1}{2}+}[0, T]$ we have $\|u(t)\|_{L^2} \leq c(T)$.*

Proof. We smoothly decompose $\widehat{u} = \widehat{u}_1 + \widehat{u}_2$ with $\text{supp } \widehat{u}_1 \subset \{|\xi| \leq 2\}$ and $\text{supp } \widehat{u}_2 \subset \{|\xi| \geq 1\}$. Then we have by Gagliardo-Nirenberg

$$\begin{aligned} \|u\|_{L^3} &\leq \|u_1\|_{L^3} + \|u_2\|_{L^3} \lesssim \|\nabla u_1\|_{L^2}^{\frac{1}{2}} \|u_1\|_{L^2}^{\frac{1}{2}} + \| |D|^{\frac{1}{2}} u_2 \|_{L^2} \\ &\lesssim \|\nabla u_1\|_{L^2}^{\frac{1}{2}} \|u_1\|_{L^2}^{\frac{1}{2}} + \| |D|^s u_2 \|_{L^2}^{\frac{1}{2s}} \|u_2\|_{L^2}^{1-\frac{1}{2s}} \\ &\lesssim \|\nabla u_1\|_{L^2}^2 + \|u_1\|_{L^2}^{\frac{2}{3}} + \|u_2\|_{L^2}^{\frac{2}{3}} + \| |D|^s u_2 \|_{L^2}^{\frac{2}{3-2s}}, \end{aligned}$$

so that by (11), (12) and (13) we get on $[0, T]$:

$$\begin{aligned} \|u(t)\|_{L^3}^3 &\lesssim \| |\xi| \widehat{u}_1(\xi, t) \|_{L^2}^6 + \|u_1(t)\|_{L^2}^2 + \|u_2(t)\|_{L^2}^2 + \| |\xi|^s \widehat{u}_2(\xi, t) \|_{L^2}^{\frac{6}{3-2s}} \\ &\leq c'(T) (\|u(t)\|_{L^2}^2 + 1). \end{aligned}$$

Multiplying the differential equation (5) with iu and taking the real part we get by Young's inequality, because $W \in L^1$ is real-valued:

$$\begin{aligned} \frac{d}{dt} \|u(t)\|_{L^2}^2 &= \int (W * (|u|^2 + 2\operatorname{Re} u)) u dx + \operatorname{Re} i \int (W * (|u|^2 + 2\operatorname{Re} u)) |u|^2 dx \\ &\lesssim \int |W * (|u|^2)| |u| dx + 2 \int |W * \operatorname{Re} u| |u| dx \\ &\lesssim \|W * |u|^2\|_{L^{\frac{3}{2}}} \|u\|_{L^3} + \|u\|_{L^2}^2 \\ &\lesssim \|u\|_{L^3}^3 + \|u\|_{L^2}^2 \\ &\leq c'(T)(\|u(t)\|_{L^2}^2 + 1) \end{aligned}$$

so that Gronwall's lemma gives

$$\|u(t)\|_{L^2}^2 + 1 \leq (\|u_0\|_{L^2}^2 + 1)e^{c'(T)T}$$

on $[0, T]$. \square

We recall our aim to give an a-priori bound of $\|\nabla Iu(t)\|_{L^2}$ (cf. (11)) on $[0, T]$ for an arbitrarily given T . We want to show this in the rest of this section as a consequence of Proposition 2.1 and the estimates for the modified energy which we give in the next section.

Let $N \geq 1$ be a number to be specified later and $s > 1/2$. Let data $u_0 \in H^s$ be given. Then we have

$$\begin{aligned} \|\nabla Iu_0\|_{L^2}^2 &\lesssim \| |\xi| \widehat{u_0}(\xi) \|_{L^2(\{|\xi| \leq N\})}^2 + \| N^{1-s} |\xi|^s \widehat{u_0}(\xi) \|_{L^2(\{|\xi| \geq N\})}^2 \\ &\lesssim \| N^{1-s} |\xi|^s \widehat{u_0}(\xi) \|_{L^2(\mathbb{R}^3)}^2 = N^{2(1-s)} \|u_0\|_{H^s}^2 \lesssim N^{2(1-s)}. \end{aligned} \quad (14)$$

This implies an estimate for $|E(Iu_0)|$ as follows: we have by Young's inequality, using $W \in L^1$:

$$\left| \int (W * (|Iu_0|^2 + 2\operatorname{Re} Iu_0)) (|Iu_0|^2 + 2\operatorname{Re} Iu_0) dx \right| \lesssim \|Iu_0\|_{L^4}^4 + \|Iu_0\|_{L^3}^3 + \|Iu_0\|_{L^2}^2.$$

Now by Sobolev's embedding

$$\begin{aligned} \|Iu_0\|_{L^4}^4 &\lesssim \|Iu_0\|_{\dot{H}^{\frac{3}{4}}}^4 \lesssim \| |\xi|^{\frac{3}{4}} \widehat{u_0} \|_{L^2(\{|\xi| \leq N\})}^4 + \left\| \frac{N^{1-s}}{|\xi|^{1-s}} |\xi|^{\frac{3}{4}} \widehat{u_0} \right\|_{L^2(\{|\xi| \geq N\})}^4 \\ &\lesssim \|\widehat{u_0}\|_{L^2(\{|\xi| \leq 1\})}^4 + \| |\xi|^{\frac{3}{4}-s} |\xi|^s \widehat{u_0} \|_{L^2(\{1 \leq |\xi| \leq N\})}^4 + N^{3-4s} \| |\xi|^s \widehat{u_0} \|_{L^2(\{|\xi| \geq N\})}^4 \\ &\lesssim \|u_0\|_{L^2}^4 + \langle N \rangle^{3-4s} \|u_0\|_{H^s}^4 \lesssim N^{2(1-s)} \|u_0\|_{H^s}^4, \end{aligned}$$

using in the last line the assumption $s \geq 1/2$. Moreover

$$\|Iu_0\|_{L^3}^3 \lesssim \|Iu_0\|_{\dot{H}^{\frac{1}{2}}}^3 \leq \|u_0\|_{\dot{H}^{\frac{1}{2}}}^3 \leq \|u_0\|_{H^s}^3$$

and

$$\|Iu_0\|_{L^2}^2 \leq \|u_0\|_{L^2}^2,$$

so that

$$|E(Iu_0)| \leq c_0 N^{2(1-s)}.$$

The local existence theorem (Prop. 2.1) implies that there exists a solution u on some time interval $[0, \delta]$ with

$$\|\nabla Iu\|_{X^{0, \frac{1}{2} + [0, \delta]}}^2 \leq M \|\nabla Iu_0\|_{L^2}^2 \leq c_0 M N^{2(1-s)+2\epsilon} \quad (15)$$

under the assumption $\|\nabla Iu_0\|_{L^2}^2 \leq c_0 N^{2(1-s+\epsilon)}$, where $\epsilon \geq 0$ and

$$\delta \sim \frac{1}{N^{2(1-s+\epsilon)}}. \quad (16)$$

Now we use the results of the next section. We have the following estimate

$$\begin{aligned} & |E(Iu(\delta)) - E(Iu_0)| \\ & \lesssim \left(\frac{\delta^{\frac{1}{2}}}{N^{2-}} + \frac{\delta^{1-}}{N^{1-}} + \frac{1}{N^{3-}} \right) \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^4 + \frac{\delta^{\frac{1}{2}}}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^3 \\ & \quad + \left(\frac{1}{N^{4-}} + \frac{\delta}{N^{2-}} \right) \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^6 + \frac{\delta}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^5. \end{aligned} \quad (17)$$

If we use (15) and (16) we easily see that the decisive term is

$$\frac{\delta^{1-}}{N^{1-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^4 + \frac{\delta}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^6 \lesssim \left(\frac{N^{4(1-s+\epsilon)}}{N^{1-}} + \frac{N^{6(1-s+\epsilon)}}{N^{2-}} \right) \delta^{1-}.$$

This is the bound for the increment of the modified energy from time 0 to time δ . Similarly we get the same bound for the increment from time $t = k\delta$ to time $t = (k+1)\delta$ for $0 \leq k \leq T/\delta$, $k \in \mathbb{N}$, provided we have a uniform bound

$$\|\nabla Iu(k\delta)\|_{L^2}^2 \leq 2c_0 N^{2(1-s+\epsilon)}, \quad (18)$$

which implies

$$\|\nabla Iu\|_{X^{0, \frac{1}{2}+}[k\delta, (k+1)\delta]}^2 \leq 2c_0 M N^{2(1-s+\epsilon)} \quad (19)$$

by the local existence theorem. The number of iteration steps to reach the given time T is T/δ , so that the increment of the energy from time $t = 0$ to time $t = (k+1)\delta$, $0 \leq k \leq \frac{T}{\delta}$, is bounded by

$$\frac{T}{\delta} \left(\frac{N^{4(1-s+\epsilon)}}{N^{1-}} + \frac{N^{6(1-s+\epsilon)}}{N^{2-}} \right) \delta^{1-} \leq c_0 N^{2(1-s)}$$

independently of k , if $TN^{1-}N^{4(1-s+\epsilon)} \ll N^{2(1-s)}$ and $TN^{-2+}N^{6(1-s+\epsilon)} \ll N^{2(1-s)}$. These conditions are fulfilled for N sufficiently large, if

$$s > \frac{1}{2} + 2\epsilon \iff \epsilon < \frac{s - \frac{1}{2}}{2}, \quad (20)$$

as one easily calculates. Choosing ϵ sufficiently small this condition is fulfilled under our assumption $s > 1/2$. We recall again that we used (18). We arrive at

$$|E(Iu(t))| \leq 2c_0 N^{2(1-s)} \quad \forall t \in [0, (k+1)\delta], \quad 0 \leq k \leq \frac{T}{\delta}, \quad k \text{ fixed}. \quad (21)$$

Now we consider the cases where either (A1) and (A2) or else (A1) and (A3) hold separately.

If (A1) and (A2) hold we have $\widehat{W}(\xi) > 0$, which immediately implies that the energy functional is positive definite, both terms in (9) are namely nonnegative, so that one gets

$$\|\nabla Iu(t)\|_{L^2}^2 \leq E(Iu(t)) \leq 2c_0 N^{2(1-s)}$$

for $0 \leq t \leq (k+1)\delta$ and $0 \leq k < \frac{T}{\delta}$, where c_0 is independent of k and where we can choose $\epsilon = 0$. Remark that on the r.h.s. the same constant $2c_0$ appears as in (18). Thus step by step after $\sim \frac{T}{\delta}$ steps we get the desired a-priori bound

$$\|\nabla Iu(t)\|_{L^2} \leq c(T) \quad \forall 0 \leq t \leq T.$$

Thus we are done in this case (modulo the results of the next section).

If (A1) and (A3) hold, the energy functional is not necessarily positive definite and it is more difficult to obtain a bound for $\|\nabla Iu(t)\|_{L^2}$ from energy bounds.

We follow the computations of de Laire [L] in this case and get

$$\begin{aligned} E(Iu) &= \|\nabla Iu\|_{L^2}^2 + \frac{1}{2} \int (W * (|Iu|^2 + 2\operatorname{Re} Iu)) (|Iu|^2 + 2\operatorname{Re} Iu) dx \\ &= \|\nabla Iu\|_{L^2}^2 + \tilde{I}_1(Iu) + \tilde{I}_2(Iu) + \tilde{I}_3(Iu), \end{aligned} \quad (22)$$

where

$$\begin{aligned}\tilde{I}_1(Iu) &:= \int (W * \operatorname{Re} Iu) \operatorname{Re} Iu dx \\ \tilde{I}_2(Iu) &:= \frac{1}{2} \int (W * |Iu|^2) (|(Iu)_1|^2 + 4 \operatorname{Re} (Iu)_1) dx \\ \tilde{I}_3(Iu) &:= \frac{1}{2} \int (W * |Iu|^2) (|(Iu)_2|^2 + 4 \operatorname{Re} (Iu)_2) dx.\end{aligned}$$

Here

$$(Iu)_1 := Iu \chi_{\{|Iu| \leq 5\}} \quad , \quad (Iu)_2 := Iu \chi_{\{|Iu| > 5\}}.$$

We used that W is even which implies

$$\int (W * |Iu|^2) \operatorname{Re} Iu dx = \int (W * \operatorname{Re} Iu) |Iu|^2 dx.$$

Using $W \in L^1(\mathbb{R}^3)$ we easily see that

$$|\tilde{I}_1(Iu)| + |\tilde{I}_2(Iu)| \lesssim \|Iu\|_{L^2}^2.$$

Moreover using (A3) we get

$$\begin{aligned}\tilde{I}_3(Iu) &\geq \frac{1}{2} \int (W * |Iu|^2) (|(Iu)_2|^2 - 4|(Iu)_2|) dx \\ &= \frac{1}{2} \int (W * |Iu|^2) |(Iu)_2| (|(Iu)_2| - 4) dx \\ &\geq \frac{1}{2} \int (W * |Iu|^2) |(Iu)_2| dx =: J_3(Iu) \geq 0.\end{aligned}\tag{23}$$

This implies by (22)

$$\|\nabla Iu\|_{L^2}^2 + \tilde{I}_3(Iu) \leq |E(Iu)| + a\|Iu\|_{L^2}^2.\tag{24}$$

In order to estimate $\|Iu\|_{L^2}^2$ we apply I to the differential equation (5), multiply with iIu and take the real part. This leads to

$$\frac{d}{dt} \|Iu\|_{L^2}^2 = \operatorname{Im} \langle F(Iu) - IF(u), Iu \rangle - \operatorname{Im} \langle F(Iu), Iu \rangle.\tag{25}$$

Let us consider the first term on the r.h.s. We get

$$\begin{aligned}\operatorname{Im} \langle F(Iu) - IF(u), Iu \rangle &= \operatorname{Im} \langle (W * |Iu|^2) Iu - I((W * |u|^2)u), Iu \rangle \\ &\quad + 2 \operatorname{Im} \langle (W * (\operatorname{Re} Iu)) Iu - I((W * \operatorname{Re} u)u), Iu \rangle \\ &\quad + \operatorname{Im} \langle (W * |Iu|^2) - I(W * |u|^2), Iu \rangle \\ &\quad + 2 \operatorname{Im} \langle W * (\operatorname{Re} Iu) - I(W * \operatorname{Re} u), Iu \rangle \\ &= \operatorname{Im} \langle I(W * |u|^2) - (W * |Iu|^2), Iu \rangle,\end{aligned}$$

Now we claim

$$\int_{k\delta}^{(k+1)\delta} \langle I(W * |u|^2) - (W * |Iu|^2), Iu \rangle dt \lesssim N^{-2} \delta \|\nabla Iu\|_{X^{0, \frac{1}{2} + [k\delta, (k+1)\delta]}}^3.\tag{26}$$

Using $|\tilde{W}(\xi)| \lesssim \frac{1}{\langle \xi \rangle^2}$ and defining

$$M(\xi_1, \xi_2) := \frac{|m(\xi_1 + \xi_2) - m(\xi_1)m(\xi_2)|}{m(\xi_1)m(\xi_2)} \cdot \frac{1}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1| |\xi_2| |\xi_3|}$$

we have to show

$$A := \int_{k\delta}^{(k+1)\delta} \int_* M(\xi_1, \xi_2) \prod_{i=1}^3 \widehat{u}_i(\xi_i, t) d\xi_1 d\xi_2 d\xi_3 dt \lesssim N^{-2} \delta \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2} +}}^3,$$

where $*$ denotes integration over $\{\sum_{i=1}^3 \xi_i = 0\}$. We assume w.l.o.g. $|\xi_1| \geq |\xi_2|$.

Case 1: $|\xi_1| \geq |\xi_2| \gtrsim N$:

We have

$$M(\xi_1, \xi_2) \lesssim \left(\frac{|\xi_1|}{N}\right)^{\frac{1}{2}} \left(\frac{|\xi_2|}{N}\right)^{\frac{1}{2}} \frac{1}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1| |\xi_2|}.$$

Thus

$$\begin{aligned} A &\lesssim N^{-2} \|u_1\|_{L_{x,t}^2} \|u_2\|_{L_{x,t}^2} \|\langle D \rangle^{-2} D^{-1} u_3\|_{L_{x,t}^\infty} \\ &\lesssim N^{-2} \delta \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2} \|D^{-1} u_3\|_{L_t^\infty L_x^6} \\ &\lesssim N^{-2} \delta \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

Case 2: $|\xi_1| \gtrsim N \gg |\xi_2|$ ($\implies |\xi_3| \sim |\xi_1| \gtrsim N$)

By the mean value theorem we get

$$M(\xi_1, \xi_2) \lesssim \frac{|\xi_2|}{|\xi_1|} \frac{1}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1| |\xi_2| |\xi_3|}.$$

Thus

$$\begin{aligned} A &\lesssim N^{-3} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^1 L_x^2} \|u_3\|_{L_t^\infty L_x^2} \lesssim N^{-3} \delta \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^\infty L_x^2} \\ &\lesssim N^{-3} \delta \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}}, \end{aligned}$$

which completes the proof of (26).

Next we estimate the last term in (25). We have

$$\begin{aligned} |Im \langle F(Iu), Iu \rangle| &= |Im \int (W * (|Iu|^2 + 2Re Iu))(1 + Iu) I \bar{u} dx| \\ &= |Im \int (W * (|Iu|^2 + 2Re Iu)) I \bar{u} dx| \\ &\leq 2 \int |W * Re Iu| |Iu| dx + \int (W * |Iu|^2) |(Iu)_1| dx + \int (W * |Iu|^2) |(Iu)_2| dx \\ &\lesssim \|Iu\|_{L^2}^2 + J_3(Iu) \lesssim \|Iu\|_{L^2}^2 + \tilde{I}_3(Iu) \lesssim \|Iu\|_{L^2}^2 + |E(Iu)| \end{aligned}$$

by (23) and (24).

From (25) we conclude for $k\delta \leq t \leq (k+1)\delta$:

$$\begin{aligned} &\|Iu(t)\|_{L^2}^2 - \|Iu(k\delta)\|_{L^2}^2 \\ &\lesssim c_1 (N^{-2} \delta \|\nabla Iu\|_{X^{0, \frac{1}{2}+[k\delta, (k+1)\delta]}}^3 + \int_{k\delta}^{(k+1)\delta} |E(Iu(s))| ds + \int_{k\delta}^t \|Iu(s)\|_{L^2}^2 ds). \end{aligned}$$

Now we have

$$c_1 (N^{-2} \delta \|\nabla Iu\|_{X^{0, \frac{1}{2}+[k\delta, (k+1)\delta]}}^3 \leq c_1 (2c_0)^{\frac{3}{2}} N^{-2} \delta N^{3(1-s)+3\epsilon} \leq c_2 \delta N^{2(1-s)},$$

provided (18) holds (and therefore (19)) and ϵ is sufficiently small. Using the uniform energy bound (21) we get for $t \in [k\delta, (k+1)\delta]$:

$$\|Iu(t)\|_{L^2}^2 \leq \|Iu(k\delta)\|_{L^2}^2 + c_2 \delta N^{2(1-s)} + 2c_0 \delta N^{2(1-s)} + c_1 \int_{k\delta}^t \|Iu(s)\|_{L^2}^2 ds.$$

Gronwall's lemma implies

$$\sup_{k\delta \leq t \leq (k+1)\delta} \|Iu(t)\|_{L^2}^2 \leq (\|Iu(k\delta)\|_{L^2}^2 + c_3 \delta N^{2(1-s)}) e^{c_1 \delta}$$

under our assumptions (21)

$$|E(Iu(t))| \leq 2c_0 N^{2(1-s)} \quad \text{on} \quad [0, (k+1)\delta]$$

and (18)

$$\|\nabla Iu(k\delta)\|_{L^2}^2 \leq 2c_0 N^{2(1-s+\epsilon)}.$$

Here c_1 and c_3 are independent of k . Using the bound for $\|Iu(k\delta)\|_{L^2}^2$ this implies

$$\begin{aligned} \sup_{k\delta \leq t \leq (k+1)\delta} \|Iu(t)\|_{L^2}^2 &\leq [(\|Iu((k-1)\delta)\|_{L^2}^2 + c_3\delta N^{2(1-s)})e^{c_1\delta} + c_3\delta N^{2(1-s)}]e^{c_1\delta} \\ &= \|Iu((k-1)\delta)\|_{L^2}^2 e^{2c_1\delta} + c_3\delta N^{2(1-s)}(e^{2c_1\delta} + e^{c_1\delta}). \end{aligned}$$

Iterating this procedure after $k \leq \frac{T}{\delta}$ steps we arrive at

$$\begin{aligned} \sup_{k\delta \leq t \leq (k+1)\delta} \|Iu(t)\|_{L^2}^2 &\leq \|Iu_0\|_{L^2}^2 e^{\frac{T}{\delta}c_1\delta} + c_3\delta N^{2(1-s)} \sum_{l=0}^{\frac{T}{\delta}} (e^{c_1\delta})^l \\ &\leq \|u_0\|_{L^2}^2 e^{c_1T} + c_4 N^{2(1-s)} e^{c_1T} \\ &\leq \frac{c_0}{a} N^{2(1-s+\epsilon)} \end{aligned}$$

choosing N so large that $e^{c_1T} \ll N^\epsilon$ with a small $\epsilon > 0$, which fulfills (20), and N also so large, that $\|u_0\|_{L^2}^2 \ll N^{2(1-s)}$. We used

$$\sum_{l=0}^{\frac{T}{\delta}} (e^{c_1\delta})^l = \frac{(e^{c_1\delta})^{\frac{T}{\delta}} - 1}{e^{c_1\delta} - 1} \lesssim \frac{e^{c_1T}}{\delta}.$$

This bound for $\|Iu(t)\|_{L^2}$ for $t \in [k\delta, (k+1)\delta]$ implies by (24),(18),(21) :

$$\begin{aligned} \|\nabla Iu(t)\|_{L^2}^2 &\leq |E(Iu(t))| + a\|Iu(t)\|_{L^2}^2 \\ &\leq 2c_0 N^{2(1-s)} + c_0 N^{2(1-s+\epsilon)} \leq 2c_0 N^{2(1-s+\epsilon)} \end{aligned}$$

for $t \in (k\delta, (k+1)\delta]$ (and choosing N so large, that $N^{2\epsilon} \geq 2$), the same bound which we had for $t = k\delta$ (cf. (18)). By iteration we thus get

$$\sup_{0 \leq t \leq T} \|\nabla Iu(t)\|_{L^2}^2 \leq 2c_0 N^{2(1-s+\epsilon)} =: c(T).$$

This completes the proof of the a-priori bound for $\|\nabla Iu(t)\|_{L^2}$ for the problem under the assumptions (A1) and (A3), so that now (11) holds in both cases. Thus the global well-posedness result is proven (modulo the results in the next section).

4. ESTIMATES FOR THE MODIFIED ENERGY

In order to estimate the increment of the modified energy $E(Iu(t))$ of a solution u of the Cauchy problem (5),(6),(7) from time t_0 to time $t_0 + \delta$, say $t_0 = 0$ for ease of notation, we have to control its time derivative. We calculate

$$\begin{aligned} \frac{d}{dt} E(Iu) &= 2\operatorname{Re} \langle -\Delta Iu, Iu_t \rangle + \frac{1}{2} \int (W * (Iu I\bar{u}_t + Iu_t I\bar{u} + 2\operatorname{Re} Iu_t))(|Iu|^2 + 2\operatorname{Re} Iu) dx \\ &\quad + \frac{1}{2} \int (W * (|Iu|^2 + 2\operatorname{Re} Iu))(Iu I\bar{u}_t + I\bar{u} Iu_t + 2\operatorname{Re} Iu_t) dx \\ &= 2\operatorname{Re} \langle -\Delta Iu, Iu_t \rangle + 2\operatorname{Re} \langle (W * (|Iu|^2 + 2\operatorname{Re} Iu))(1 + Iu), Iu_t \rangle, \end{aligned}$$

where we used that W is even, so that the second and third term coincide. Now

$$-\Delta Iu = -iIu_t - IF(u),$$

so that

$$\begin{aligned} \frac{d}{dt}E(Iu) &= 2\operatorname{Re}\langle F(Iu) - IF(u), Iu_t \rangle \\ &= 2\operatorname{Im}(\langle \nabla(F(Iu) - IF(u)), \nabla Iu \rangle - \langle F(Iu) - IF(u), IF(u) \rangle) \end{aligned}$$

and

$$|\frac{d}{dt}E(Iu)| \leq 2(|\langle \nabla(F(Iu) - IF(u)), \nabla Iu \rangle| + |\langle F(Iu) - IF(u), IF(u) \rangle|) \quad (27)$$

with (cf. (7))

$$F(u) = (1 + u)(W * (|u|^2 + 2\operatorname{Re} u)).$$

This especially shows the standard energy conservation law (setting $I = id$).

The estimates which now follow are given in terms of bounds of Fourier transforms of the corresponding functions. The only property of W which we use is the bound $|\widehat{W}(\xi)| \lesssim \langle \xi \rangle^{-2}$, so that both cases, namely assuming either (A1) and (A2) or else (A1) and (A3) can be handled in the same way. The most critical cases are the terms of fourth and third order of the first term on the r.h.s. of (27) and the term of sixth order of the second term. In fact we shall refer to the estimates in the case of the local Gross-Pitaevskii equation, where $\widehat{W} = 1$, in our earlier paper [Pe] for the remaining terms of lower order on the r.h.s. of (27).

We start with the terms of highest order in the first term. Taking again the time interval $[0, \delta]$ instead of $[k\delta, (k+1)\delta]$ just for the ease of notation we claim

$$|\int_0^\delta \langle \nabla((W * |Iu|^2)Iu - I((W * |u|^2)u)), \nabla Iu \rangle dt| \lesssim (\frac{\delta^{1-}}{N^{1-}} + \frac{\delta^{\frac{1}{2}}}{N^{2-}}) \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^4. \quad (28)$$

Here and in the following we use dyadic decompositions with respect to the space variables ξ_i , where $|\xi_i| \sim N_i$ with $N_i = 2^{k_i}$, $k_i \in \mathbb{Z}$. In order to sum the dyadic parts at the end we always need a convergence generating factor $\frac{1 \wedge N_{min}^{0+}}{N_{max}^{0+}}$, where N_{min} and N_{max} is the smallest and the largest of the numbers N_i , respectively. $N_{max} \geq N(\geq 1)$ can be assumed in all cases, because otherwise our multiplier M is identically zero. We have to take care of low frequencies especially, because we need an estimate in terms of ∇Iu . Assuming w.l.o.g. that the Fourier transforms are nonnegative we have to show:

$$A := \int_0^\delta \int_* M(\xi_1, \xi_2, \xi_3) \prod_{i=1}^4 \widehat{u}_i(\xi, t) d\xi_1 \dots d\xi_4 dt \lesssim (\frac{\delta^{1-}}{N^{1-}} + \frac{\delta^{\frac{1}{2}}}{N^{2-}}) \prod_{i=1}^4 \|u_i\|_{X^{0, \frac{1}{2}+}[0, \delta]},$$

where $*$ always denotes integration over $\{\sum \xi_i = 0\}$, and

$$M(\xi_1, \xi_2, \xi_3) := \frac{|m(\xi_1 + \xi_2 + \xi_3) - m(\xi_1)m(\xi_2)m(\xi_3)|}{m(\xi_1)m(\xi_2)m(\xi_3)} \cdot \frac{|\xi_1 + \xi_2 + \xi_3|}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1| |\xi_2| |\xi_3|}.$$

Case 1: $N_3 \gg N_1 \geq N_2$.

In this case we have $N_3 \sim N_4 \gtrsim N$, $N_2 = N_{min}$ and $N_3 \sim N_{max}$. By the mean value theorem we get

$$M(\xi_1, \xi_2, \xi_3) \lesssim \frac{N_1}{N_3} \cdot \frac{1}{N_1 |\xi_2| \langle \xi_1 + \xi_2 \rangle^2}$$

and by Hölder's inequality, Sobolev's embedding and Strichartz' estimate we get

$$\begin{aligned}
A &\lesssim \frac{1}{N_3} \|\langle D \rangle^{-2} (D^{-1} u_2 u_1)\|_{L_t^\infty L_x^\infty} \|u_3\|_{L_t^{2+} L_x^{2+}} \|u_4\|_{L_t^2 L_x^2} \\
&\lesssim \frac{\delta^{1-}}{N_3} \|D^{-1} u_2 u_1\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^{2+}} \|u_4\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-}}{N_3} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_1\|_{L_t^\infty L_x^{2+}} \|u_3\|_{L_t^\infty L_x^{2+}} \|u_4\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-} (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{1-}} \prod_{i=1}^4 \|u_i\|_{X^{0, \frac{1}{2}+}},
\end{aligned}$$

using $\|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \lesssim N_2^{0+} \|u_2\|_{L_t^\infty L_x^2}$ and $\|u_3\|_{L_t^\infty L_x^{2+}} \lesssim N_3^{0+} \|u_3\|_{L_t^\infty L_x^2}$.

Case 2: $N_1 \gg N_2, N_3$ and $N_1 \gtrsim N$.

We have similarly as in case 1:

$$M(\xi_1, \xi_2, \xi_3) \lesssim \frac{N_3}{N_1} \cdot \frac{1}{|\xi_2| N_3 \langle \xi_1 + \xi_2 \rangle^2}$$

and get the same estimate as in case 1 interchanging the roles of N_1 and N_3 .

Case 3: $N_1 \sim N_3$ ($\implies N_1, N_3 \gtrsim N$)

In this case we get

$$\begin{aligned}
M(\xi_1, \xi_2, \xi_3) &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \left(\frac{N_3}{N}\right)^{\frac{1}{2}-} \left\langle \frac{N_1}{N} \right\rangle^{\frac{1}{2}-} \frac{1}{|\xi_2| |\xi_3| \langle \xi_1 + \xi_2 \rangle^2} \\
&\lesssim \frac{1}{N_3^{0+} N^{1-}} \left\langle \frac{N_2}{N} \right\rangle^{\frac{1}{2}-} \frac{1}{N_2 \langle \xi_1 + \xi_2 \rangle^2}.
\end{aligned}$$

a. $N_2 \gtrsim N$.

We have

$$\begin{aligned}
A &\lesssim \frac{1}{N_3^{0+} N^{2-}} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^2 L_x^{3-}} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^2 L_x^{6+}} \\
&\lesssim \frac{\delta^{\frac{1}{2}} N_4^{0+}}{N_3^{0+} N^{2-}} \|u_1 u_2\|_{L_t^\infty L_x^{1+}} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{X^{0, \frac{1}{2}+}} \\
&\lesssim \frac{\delta^{\frac{1}{2}} (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^4 \|u_i\|_{X^{0, \frac{1}{2}+}},
\end{aligned}$$

using $\|u_4\|_{L_t^2 L_x^{6+}} \lesssim N_4^{0+} \|u_4\|_{L_t^2 L_x^6} \lesssim \|u_4\|_{X^{0, \frac{1}{2}+}}$ by Strichartz' estimate.

b. $N_2 \ll N$.

In this case we get

$$\begin{aligned}
A &\lesssim \frac{1}{N_{\max}^{0+} N^{1-}} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^1 L_x^\infty} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty L_x^{2+}} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{1-}} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty L_x^{2+}} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{1-}} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty L_x^{2+}} \\
&\lesssim \frac{\delta (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{1-}} \prod_{i=1}^4 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

Case 4: $N_1 \sim N_2 \gtrsim N_3$ ($\implies N_1, N_2 \gtrsim N$)

In this case we get

$$\begin{aligned} M(\xi_1, \xi_2, \xi_3) &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \left(\frac{N_2}{N}\right)^{\frac{1}{2}-} \left\langle \frac{N_3}{N} \right\rangle^{\frac{1}{2}-} \frac{1}{|\xi_1||\xi_3|\langle \xi_1 + \xi_2 \rangle^2} \\ &\lesssim \frac{1}{N_{\max}^{0+} N^{1-}} \left\langle \frac{N_3}{N} \right\rangle^{\frac{1}{2}-} \frac{1}{N_3 \langle \xi_1 + \xi_2 \rangle^2}. \end{aligned}$$

The case $N_3 \gtrsim N$ is handled like case 3a, whereas in the case $N_3 \ll N$ we get

$$\begin{aligned} A &\lesssim \frac{1}{N_{\max}^{0+} N^{1-}} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^1 L_x^{3-}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \|u_4\|_{L_t^\infty L_x^{2+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{1-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{1-}} \prod_{i=1}^4 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

Case 5: $N_2 \sim N_3$ ($\implies N_1 \gtrsim N_2 \sim N_3$).

If $N_1 \gg N_2$ we are in the situation of case 2, otherwise $N_1 \sim N_2 \sim N_3 \gtrsim N$, so that

$$\begin{aligned} M(\xi_1, \xi_2, \xi_3) &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \left(\frac{N_2}{N}\right)^{\frac{1}{2}-} \left(\frac{N_3}{N}\right)^{\frac{1}{2}-} \frac{1}{N_2 N_3 \langle \xi_1 + \xi_2 \rangle^2} \\ &\lesssim \frac{1}{N_{\max}^{0+} N^{2-}} \frac{1}{N_2 \langle \xi_1 + \xi_2 \rangle^2} \end{aligned}$$

as in case 3a.

Dyadic summation gives estimate (28).

We next consider the cubic part of the first term on the r.h.s. of (27).

We claim

$$\left| \int_0^\delta \langle \nabla(W * (|Iu|^2) - I(W * |u|^2)), \nabla Iu \rangle dt \right| \lesssim \frac{\delta^{\frac{1}{2}}}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+[0, \delta]}}^3. \quad (29)$$

We have to show

$$B := \int_0^\delta \int_* M(\xi_1, \xi_2) \prod_{i=1}^3 \widehat{u}_i(\xi_i, t) d\xi_1 d\xi_2 d\xi_3 dt \lesssim \frac{\delta^{\frac{1}{2}}}{N^{2-}} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+[0, \delta]}}$$

with

$$M(\xi_1, \xi_2) := \frac{|m(\xi_1 + \xi_2) - m(\xi_1)m(\xi_2)|}{m(\xi_1)m(\xi_2)} \cdot \frac{|\xi_1 + \xi_2|}{\langle \xi_1 + \xi_2 \rangle^2 |\xi_1||\xi_2|}.$$

Case 1: $N_1 \geq N_2 \gtrsim N$.

We have

$$M(\xi_1, \xi_2) \lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_2}{N}\right)^{\frac{1}{2}} \frac{1}{N_1 N_2 \langle \xi_1 + \xi_2 \rangle},$$

so that by Strichartz' estimate:

$$\begin{aligned} B &\lesssim \frac{1}{N_{\max}^{0+} N^{2-}} \|\langle D \rangle^{-1} (u_1 u_2)\|_{L_t^1 L_x^{2-}} \|u_3\|_{L_t^\infty L_x^{2+}} \\ &\lesssim \frac{1}{N_{\max}^{0+} N^{2-}} \|u_1 u_2\|_{L_t^1 L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^{2+}} \\ &\lesssim \frac{\delta^{\frac{1}{2}}}{N_{\max}^{0+} N^{2-}} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^2 L_x^{6+}} \|u_3\|_{L_t^\infty L_x^{2+}} \\ &\lesssim \frac{\delta^{\frac{1}{2}}(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

Case 2: $N_1 \geq N \gg N_2$.

Similarly we get by the mean value theorem

$$M(\xi_1, \xi_2) \lesssim \frac{N_2}{N_1} \frac{1}{N_1 N_2 \langle \xi_1 + \xi_2 \rangle} \lesssim \frac{1}{N_{max}^{0+} N^{2-} \langle \xi_1 + \xi_2 \rangle}$$

leading to the same bound as in case 1, thus (29) is proven.

Concerning the second cubic term we claim

$$|\int_0^\delta \langle \nabla((W * Iu)Iu - I((W * u)u)), \nabla Iu \rangle dt| \lesssim \frac{\delta}{N^{1-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^3. \quad (30)$$

We again have to consider a term like B but with

$$M(\xi_1, \xi_2) := \frac{|m(\xi_1 + \xi_2) - m(\xi_1)m(\xi_2)|}{m(\xi_1)m(\xi_2)} \cdot \frac{|\xi_1 + \xi_2|}{\langle \xi_1 \rangle^2 |\xi_1| |\xi_2|}.$$

We concentrate on the more difficult case $N_2 \geq N_1$ and have to consider

Case 1: $N_2 \sim N_1 \gtrsim N$.

We have

$$M(\xi_1, \xi_2) \lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \left(\frac{N_2}{N}\right)^{\frac{1}{2}-} \frac{1}{N_1^3} \lesssim \frac{1}{N_{max}^{0+} N_1^{3-}}.$$

Thus

$$\begin{aligned} B &\lesssim \frac{\delta}{N_{max}^{0+} N_1^{3-}} \|u_1\|_{L_t^\infty L_x^\infty} \|u_2\|_{L_t^\infty L_x^2} \|u_3\|_{L_t^\infty L_x^{2+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{max}^{0+} N_1^{3-}} N_1^{\frac{3}{2}} \prod_{i=1}^3 \|u_i\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{max}^{0+} N^{\frac{3}{2}-}} \prod_{i=1}^3 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

Case 2: $N_2 \gg N_1 (\implies N_2 \gtrsim N)$.

By the mean value theorem we get

$$M(\xi_1, \xi_2) \lesssim \frac{N_1}{N_2} \frac{1}{\langle \xi_1 \rangle^2 N_1} = \frac{1}{N_2 \langle \xi_1 \rangle^2},$$

so that

$$\begin{aligned} B &\lesssim \frac{\delta}{N_{max}^{0+} N^{1-}} \|\langle D \rangle^{-2} u_1\|_{L_t^\infty L_x^\infty} \|u_2\|_{L_t^\infty L_x^2} \|u_3\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{max}^{0+} N^{1-}} \prod_{i=1}^3 \|u_i\|_{L_t^\infty L_x^2}. \end{aligned}$$

Thus (30) follows.

We now have to consider the sixth order term on the r.h.s. of (27). Our aim is to show the following estimate:

$$|\int_0^\delta \langle (W * |Iu|^2)Iu - I((W * |u|^2)u), I((W * |u|^2)u) \rangle dt| \lesssim \left(\frac{\delta}{N^{2-}} + \frac{1}{N^{4-}}\right) \|\nabla Iu\|_{X^{0, \frac{1}{2}+}}^6. \quad (31)$$

We have to show

$$C := \int_0^\delta \int_* M(\xi_1, \dots, \xi_6) \prod_{i=1}^6 \widehat{u}_i(\xi_i, t) d\xi_1 \dots d\xi_6 dt \lesssim \left(\frac{\delta}{N^{2-}} + \frac{1}{N^{4-}}\right) \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}[0, \delta]}$$

with

$$M(\xi_1, \dots, \xi_6) := \frac{|m(\xi_1 + \xi_2 + \xi_3) - m(\xi_1)m(\xi_2)m(\xi_3)|}{m(\xi_1)m(\xi_2)m(\xi_3)\langle \xi_1 + \xi_2 \rangle^2} \cdot \frac{m(\xi_4 + \xi_5 + \xi_6)}{m(\xi_4)m(\xi_5)m(\xi_6)\langle \xi_4 + \xi_5 \rangle^2} \cdot \prod_{i=1}^6 |\xi_i|^{-1}.$$

We assume w.l.o.g. $N_1 \geq N_2$ and $N_4 \geq N_5$.

Case 1: $N \gg N_4 \geq N_5$ and $N \gg N_6$.

a. $N_1 \geq N_2 \gtrsim N \gtrsim N_3$.

In this case we get

$$\begin{aligned} C &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_2}{N}\right)^{\frac{1}{2}} \frac{1}{N_1 N_2} \|\langle D \rangle^{-2}(u_1 u_2)\|_{L_t^1 L_x^{\frac{3}{2}-}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|\langle D \rangle^{-2}(D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_4\|_{L_t^\infty L_x^6} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

b. $N_1, N_3 \gtrsim N \gg N_2$.

The estimate

$$M(\xi_1, \dots, \xi_6) \lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_3}{N}\right)^{\frac{1}{2}} \frac{1}{N_1^2 \langle \xi_4 + \xi_5 \rangle^2 |\xi_1| |\xi_2| N_3 |\xi_4| |\xi_5| |\xi_6|}$$

implies

$$\begin{aligned} C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|D^{-1} u_1\|_{L_t^\infty L_x^6} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|\langle D \rangle^{-2}(D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \\ &\quad \cdot \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

c. $N_1, N_2, N_3 \gtrsim N$.

In this case we get

$$\begin{aligned} C &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_2}{N}\right)^{\frac{1}{2}} \left(\frac{N_3}{N}\right)^{\frac{1}{2}} \frac{1}{N_1 N_2 N_3} \|\langle D \rangle^{-2}(u_1 u_2)\|_{L_t^1 L_x^{3-}} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|\langle D \rangle^{-2}(D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

Case 2: $N \gg N_4 \geq N_5$ and $N_6 \gtrsim N$.

a. $N_3 \gtrsim N \gtrsim N_1 \geq N_2$.

By the mean value theorem we have

$$\begin{aligned}
C &\lesssim \frac{N_1}{N_3} \cdot \frac{\delta^{1-}}{N_1 N_3 N_6} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^\infty} \|u_3\|_{L_t^\infty L_x^2} \\
&\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-}}{N_{\max}^{0+} N^{3-}} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-}}{N_{\max}^{0+} N^{3-}} \|u_1\|_{L_t^\infty L_x^{2+}} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4\|_{L_t^\infty L_x^6} \\
&\quad \cdot \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-} (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_1 \gtrsim N \gg N_2 \geq N_3$ (or $N_1 \gtrsim N \gg N_3 \geq N_2$ by exchanging u_2 and u_3).

Similarly as in a. we use the mean value theorem and get

$$\begin{aligned}
C &\lesssim \frac{N_2}{N_1} \cdot \frac{\delta}{N_1^2 N_2} \|D^{-1} u_1\|_{L_t^\infty L_x^6} \|u_2\|_{L_t^\infty L_x^2} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \|D^{-1} u_6\|_{L_t^\infty L_x^6} \\
&\lesssim \frac{\delta}{N_1^3} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

c. $N_1 \geq N_2 \gtrsim N \gg N_3$.

We get

$$\begin{aligned}
C &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \left(\frac{N_2}{N}\right)^{\frac{1}{2}-} \frac{1}{N_1 N_2 N_6} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^2 L_x^3} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^2 L_x^\infty} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

d. $N_1 \geq N_2 \geq N_3 \gtrsim N$ or $N_1 \geq N_3 \geq N_2 \gtrsim N$.

This case can be handled similarly as case c. with an additional factor $(\frac{N_3}{N})^{\frac{1}{2}-}$.

e. $N_3 \gtrsim N_1 \gtrsim N \gtrsim N_2$ (or $N_1 \geq N_3 \gtrsim N \gtrsim N_2$ by exchanging the roles of u_1 and u_3).

We get

$$\begin{aligned}
C &\lesssim \left(\frac{N_3}{N}\right)^{\frac{1}{2}-} \left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \frac{\delta}{N_1 N_3 N_6} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^\infty} \|u_3\|_{L_t^\infty L_x^2} \\
&\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4\|_{L_t^\infty L_x^6} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

f. $N_3 \geq N_1 \geq N_2 \gtrsim N$ (or $N_1 \geq N_3 \geq N_2 \gtrsim N$).

This case can be treated as case a. with an additional factor $(\frac{N_2}{N})^{\frac{1}{2}-}$.

Case 3: $N_4 \geq N_5 \gtrsim N$.

a. $N_1, N_2, N_3 \lesssim N$ and $N_6 \leq N$.

We get

$$M(\xi_1, \dots, \xi_6) \lesssim \frac{1}{\langle \xi_1 + \xi_2 \rangle^2 \langle \xi_4 + \xi_5 \rangle^2 N_4 N_5 |\xi_1| |\xi_2| |\xi_3| |\xi_6|} \left(\frac{N_4}{N}\right)^{\frac{1}{2}-} \left(\frac{N_5}{N}\right)^{\frac{1}{2}-},$$

so that

$$\begin{aligned}
C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|\langle D \rangle^{-2} (D^{-1} u_1 D^{-1} u_2)\|_{L_t^\infty L_x^\infty} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^{\frac{3}{2}}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^{3+}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \|u_4 u_5\|_{L_t^\infty L_x^1} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_1, N_2, N_3 \lesssim N$ and $N_6 \geq N$.

We argue similarly as in case a with an additional factor $(\frac{N_6}{N})^{\frac{1}{2}-}$ and get

$$\begin{aligned}
C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|\langle D \rangle^{-2} (D^{-1} u_1 D^{-1} u_2)\|_{L_t^\infty L_x^\infty} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^3} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

c. $N_3 \gtrsim N$ and $N_1, N_2, N_6 \lesssim N$.

Replacing $\|D^{-1} u_3\|_{L_t^\infty L_x^{6+}}$ by $\|D^{-1} u_3\|_{L_t^\infty L_x^6}$ we get the same result as in case a.

d. $N_3 \gtrsim N$, $N_1, N_2 \lesssim N$ and $N_6 \gtrsim N$.

The additional factor $(\frac{N_6}{N})^{\frac{1}{2}}$ leads to the bound

$$\begin{aligned}
C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|\langle D \rangle^{-2} (D^{-1} u_1 D^{-1} u_2)\|_{L_t^\infty L_x^\infty} \|D^{-1} u_3\|_{L_t^\infty L_x^6} \\
&\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^3} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

e. $N_1 \geq N_2 \gtrsim N$ and $N_3, N_6 \lesssim N$.

We get

$$M(\xi_1, \dots, \xi_6) \lesssim \frac{(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_2}{N})^{\frac{1}{2}-} (\frac{N_4}{N})^{\frac{1}{2}-} (\frac{N_5}{N})^{\frac{1}{2}-}}{\langle \xi_1 + \xi_2 \rangle^2 \langle \xi_4 + \xi_5 \rangle^2 N_1 N_2 N_4 N_5 |\xi_3| |\xi_6|},$$

so that

$$\begin{aligned} C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^{3-}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^3} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^\infty L_x^2} \|u_4 u_5\|_{L_t^\infty L_x^1} \|u_6\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

f. $N_1 \geq N_2 \gtrsim N$, $N_3 \lesssim N$ and $N_6 \gtrsim N$.

The additional factor $(\frac{N_6}{N})^{\frac{1}{2}}$ leads to

$$\begin{aligned} C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^{3-}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|\langle D \rangle^{-2} (u_4 D^{-1} u_5)\|_{L_t^\infty L_x^{\infty-}} \|u_6\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty L_x^2} \|D^{-1} u_5\|_{L_t^\infty L_x^6} \|u_6\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

g. $N_1 \geq N_2 \gtrsim N$, $N_6 \lesssim N$ and $N_3 \gtrsim N$.

This case can be treated as case f. with u_3 and u_6 exchanged.

h. $N_1 \geq N_2 \gtrsim N$ and $N_3, N_6 \gtrsim N$ ($\implies N_{\min} \gtrsim N$).

We get

$$M(\xi_1, \dots, \xi_6) \lesssim \prod_{i=1}^6 (\frac{N_i}{N})^{\frac{1}{2}-} \prod_{i=1}^6 N_i^{-1} \frac{1}{\langle \xi_1 + \xi_2 \rangle^2 \langle \xi_4 + \xi_5 \rangle^2},$$

so that

$$\begin{aligned} C &\lesssim \frac{1}{N_{\max}^{0+} N^{6-}} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^3} \|u_3\|_{L_t^2 L_x^6} \\ &\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^3} \|u_6\|_{L_t^2 L_x^6} \\ &\lesssim \frac{1}{N_{\max}^{0+} N^{6-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{X^{0, \frac{1}{2}+}} \|u_4 u_5\|_{L_t^\infty L_x^1} \|u_6\|_{X^{0, \frac{1}{2}+}} \\ &\lesssim \frac{1}{N_{\max}^{0+} N^{6-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

i. $N_1 \gtrsim N \gg N_2$ and $N_3, N_6 \lesssim N$.

We get

$$M(\xi_1, \dots, \xi_6) \lesssim \frac{(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_2}{N})^{\frac{1}{2}-} (\frac{N_4}{N})^{\frac{1}{2}-} (\frac{N_5}{N})^{\frac{1}{2}-}}{N_1^2 \langle \xi_4 + \xi_5 \rangle^2 N_1 N_2 |\xi_3| N_4 N_5 |\xi_6|},$$

so that

$$\begin{aligned} C &\lesssim \frac{1}{N_{\max}^{0+} N^{6-}} \|u_1\|_{L_t^2 L_x^6} \|u_2\|_{L_t^2 L_x^{6+}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^{3-}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{1 \wedge N_{\min}^{0+}}{N_{\max}^{0+} N^{6-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

j. $N_1 \gtrsim N \gg N_2$ and $N_3, N_6 \gtrsim N$.

We get

$$M(\xi_1, \dots, \xi_6) \lesssim \frac{(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_3}{N})^{\frac{1}{2}-} (\frac{N_4}{N})^{\frac{1}{2}-} (\frac{N_5}{N})^{\frac{1}{2}-} (\frac{N_6}{N})^{\frac{1}{2}-}}{N_1^2 \langle \xi_4 + \xi_5 \rangle^2 |\xi_1| |\xi_2| N_3 N_4 N_5 N_6},$$

thus using $N_1 \lesssim \max(N_3, N_4, N_5, N_6)$:

$$\begin{aligned} C &\lesssim \frac{1}{N_{\max}^{0+} N^{6-}} \|D^{-1} u_1\|_{L_t^\infty L_x^6} \|D^{-1} u_2\|_{L_t^2 L_x^{6+}} \|u_3\|_{L_t^2 L_x^6} \\ &\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^{3-}} \|u_6\|_{L_t^2 L_x^{6+}} \\ &\lesssim \frac{1 \wedge N_{\min}^{0+}}{N_{\max}^{0+} N^{6-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

k. $N_1 \gtrsim N \gg N_2$, $N_3 \lesssim N$ and $N_6 \lesssim N$.

This case can be handled like j. without the factor $(\frac{N_6}{N})^{\frac{1}{2}-}$ by exchanging u_1 and u_6 .

l. $N_1 \gtrsim N \gg N_2$, $N_3 \lesssim N$ and $N_6 \gtrsim N$.

The mean value theorem gives the bound

$$\begin{aligned} C &\lesssim \frac{N_2 (\frac{N_4}{N})^{\frac{1}{2}-} (\frac{N_5}{N})^{\frac{1}{2}-} (\frac{N_6}{N})^{\frac{1}{2}-}}{N_1^2 N_2 N_4 N_5 N_6} \|D^{-1} u_1\|_{L_t^\infty L_x^6} \|u_2\|_{L_t^2 L_x^{6+}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|\langle D \rangle^{-2} (u_4 u_5)\|_{L_t^\infty L_x^{3-}} \|u_6\|_{L_t^2 L_x^6} \\ &\lesssim \frac{1 \wedge N_{\min}^{0+}}{N_{\max}^{0+} N^{6-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

Case 4: $N_4 \geq N \gg N_5, N_6$.

a. $N_3 \gtrsim N \gtrsim N_1 \geq N_2$.

The mean value theorem gives

$$\begin{aligned} C &\lesssim \frac{N_1}{N_3 N_1 N_3 N_4} \delta \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^6} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|\langle D \rangle^{-2} (u_4 D^{-1} u_5)\|_{L_t^\infty L_x^{6-}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|u_4 D^{-1} u_5\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|u_4\|_{L_t^\infty L_x^2} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

b. $N_1 \geq N \gg N_2 \geq N_3$.

The mean value theorem implies

$$\begin{aligned} C &\lesssim \frac{N_2}{N_1} \cdot \frac{1}{N_1^2 N_1 N_2} \|u_1\|_{L_t^2 L_x^6} \|u_2\|_{L_t^2 L_x^6} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^{3-}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{1 \wedge N_{\min}^{0+}}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

c. $N_1 \geq N \gg N_3 \geq N_2$.

This case can be treated like case b. with u_2 and u_3 exchanged.

d. $N_3 \geq N_1 \gtrsim N \gg N_2$.

We get

$$\begin{aligned} C &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_3}{N}\right)^{\frac{1}{2}} \frac{1}{N_1^2 N_1 N_3} \|u_1\|_{L_t^2 L_x^6} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^2 L_x^6} \\ &\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^{3-}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{1 \wedge N_{\min}^{0+}}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

e. $N_1 \geq N_2 \gtrsim N \gtrsim N_3$.

We get in this case

$$M(\xi_1, \dots, \xi_6) \lesssim \frac{\left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \left(\frac{N_2}{N}\right)^{\frac{1}{2}-}}{\langle \xi_1 + \xi_2 \rangle^2 N_4^2 N_1 N_2 |\xi_3| |\xi_4| |\xi_5| |\xi_6|},$$

so that

$$\begin{aligned} C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^{3-}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\ &\quad \cdot \|D^{-1} u_4\|_{L_t^\infty L_x^6} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \prod_{i=3}^6 \|u_i\|_{L_t^\infty L_x^2} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

f. $N_1 \geq N_3 \gtrsim N \gtrsim N_2$.

We get

$$\begin{aligned} C &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}-} \left(\frac{N_3}{N}\right)^{\frac{1}{2}-} \frac{\delta}{N_4^2 N_1 N_3} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^{\infty-}} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|D^{-1} u_4\|_{L_t^\infty L_x^6} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|D^{-1} u_4\|_{L_t^\infty L_x^6} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \\ &\quad \cdot \|D^{-1} u_4\|_{L_t^\infty L_x^6} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_6\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

g. $N_1 \geq N_2 \gtrsim N$ and $N_3 \gtrsim N$.

This case can be treated as case f. with an additional factor $(\frac{N_2}{N})^{\frac{1}{2}-}$.

Case 5: $N_4, N_6 \geq N \gg N_5$.

a. $N_3 \gtrsim N \gtrsim N_1 \geq N_2$.

This case can be treated as case 2a, because the additional factor $(\frac{N_4}{N})^{\frac{1}{2}-}(\frac{N_6}{N})^{\frac{1}{2}-}$ is harmless, when one uses $N_4 \lesssim \max(N_3, N_6)$.

b. $N_1 \gtrsim N \gg N_2 \geq N_3$ (or $N_1 \gtrsim N \gg N_3 \geq N_2$ by exchanging u_2 and u_3).

We have by the mean value theorem

$$\begin{aligned}
C &\lesssim \frac{N_2}{N_1} \left(\frac{N_4}{N}\right)^{\frac{1}{2}-} \left(\frac{N_6}{N}\right)^{\frac{1}{2}-} \frac{1}{N_1^2 N_1 N_2 N_4 N_6} \|u_1\|_{L_t^\infty L_x^{2+}} \|u_2\|_{L_t^2 L_x^6} \|D^{-1}u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|\langle D \rangle^{-2}(u_4 D^{-1}u_5)\|_{L_t^\infty L_x^\infty} \|u_6\|_{L_t^2 L_x^6} \\
&\lesssim \frac{1}{N_{\max}^{0+} N^{6-}} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{X^{0, \frac{1}{2}+}} \|D^{-1}u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|u_4 D^{-1}u_5\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_6\|_{X^{0, \frac{1}{2}+}} \\
&\lesssim \frac{1}{N_{\max}^{0+} N^{6-}} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{X^{0, \frac{1}{2}+}} \|D^{-1}u_3\|_{L_t^\infty L_x^{6+}} \|u_4\|_{L_t^\infty L_x^2} \|D^{-1}u_5\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|u_6\|_{X^{0, \frac{1}{2}+}} \\
&\lesssim \frac{1 \wedge N_{\min}^{0+}}{N_{\max}^{0+} N^{6-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

c. $N_1 \geq N_2 \gtrsim N \gg N_3$.

This case is treated like case 2c, because the additional factor $(\frac{N_4}{N})^{\frac{1}{2}-}(\frac{N_6}{N})^{\frac{1}{2}-}$ can be handled using $N_4 \lesssim \max(N_1, N_6)$.

d. $N_3 \geq N_1 \gtrsim N \gg N_2$ or $N_1 \geq N_3 \gtrsim N \geq N_2$.

This case can be handled like case 2e, using $N_4 \lesssim \max(N_1, N_3, N_6)$.

e. $N_1 \geq N_2 \gtrsim N$ and $N_3 \gtrsim N$.

This case is also treated like case 2e, because the additional factor $(\frac{N_2}{N})^{\frac{1}{2}-}(\frac{N_4}{N})^{\frac{1}{2}-}(\frac{N_6}{N})^{\frac{1}{2}-}$ is acceptable, using $N_2 \leq N_1$ and $N_4 \lesssim \max(N_1, N_3, N_6)$.

Case 6: $N_4, N_5, N_6 \gtrsim N$.

a. $N_3 \gtrsim N \gg N_1, N_2$.

The mean value theorem allows to estimate

$$M(\xi_1, \dots, \xi_6) \lesssim \frac{N_1}{N_3} \cdot \frac{(\frac{N_4}{N})^{\frac{1}{2}-}(\frac{N_5}{N})^{\frac{1}{2}-}(\frac{N_6}{N})^{\frac{1}{2}-}}{\langle \xi_1 + \xi_2 \rangle^2 \langle \xi_4 + \xi_5 \rangle^2 N_1 |\xi_2| |\xi_3| N_4 N_5 N_6} \lesssim \frac{1}{N_{\max}^{0+} N^{4-}},$$

thus

$$\begin{aligned}
C &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|\langle D \rangle^{-2}(u_1 D^{-1}u_2)\|_{L_t^\infty L_x^\infty} \|D^{-1}u_3\|_{L_t^\infty L_x^6} \\
&\quad \cdot \|\langle D \rangle^{-2}(u_4 u_5)\|_{L_t^\infty L_x^3} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|u_1 D^{-1}u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \|u_4 u_5\|_{L_t^\infty L_x^1} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|u_1\|_{L_t^\infty L_x^2} \|D^{-1}u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty L_x^2} \|u_5\|_{L_t^\infty L_x^2} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_1, N_2, N_3 \gtrsim N$ and w.l.o.g. $N_1 \geq N_2$ and $N_4 \geq N_5$.

We have

$$\begin{aligned}
C &\lesssim \prod_{i=1}^6 \left(\frac{N_i}{N}\right)^{\frac{1}{2}-} \frac{\delta^{1-}}{N_1 N_3 N_4 N_6} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty - L_x^\infty} \|u_3\|_{L_t^\infty L_x^2} \\
&\quad \cdot \|\langle D \rangle^{-2} (u_4 D^{-1} u_5)\|_{L_t^\infty - L_x^\infty} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-}}{N_{\max}^{0+} N^{4-}} \|u_1 D^{-1} u_2\|_{L_t^\infty - L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \|u_4 D^{-1} u_5\|_{L_t^\infty - L_x^{\frac{3}{2}+}} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-}}{N_{\max}^{0+} N^{4-}} \|u_1\|_{L_t^\infty - L_x^{2+}} \|D^{-1} u_2\|_{L_t^\infty L_x^{6+}} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty - L_x^{2+}} \|D^{-1} u_5\|_{L_t^\infty L_x^6} \\
&\quad \cdot \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta^{1-} (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

c. $N_1, N_3 \gtrsim N \gtrsim N_2$.

This case can be treated as case b. without the factor $(\frac{N_2}{N})^{\frac{1}{2}-}$.

d. $N_1 \geq N_2 \gtrsim N \gtrsim N_3$ and w.l.o.g. $N_4 \geq N_5$.

We get

$$\begin{aligned}
C &\lesssim \prod_{i \neq 3} \left(\frac{N_i}{N}\right)^{\frac{1}{2}-} \frac{\delta}{N_1 N_2 N_4 N_6} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^3} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|\langle D \rangle^{-2} (u_4 D^{-1} u_5)\|_{L_t^\infty L_x^6} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|u_3\|_{L_t^\infty L_x^2} \|u_4 D^{-1} u_5\|_{L_t^\infty L_x^{\frac{3}{2}}} \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2} \|u_3\|_{L_t^\infty L_x^2} \|u_4\|_{L_t^\infty L_x^2} \|D^{-1} u_5\|_{L_t^\infty L_x^6} \\
&\quad \cdot \|u_6\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta (1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^6 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

This completes the proof of (31).

We now start to consider the fifth order terms and claim

$$\left| \int_0^\delta \langle (W * |Iu|^2) Iu - I((W * |u|^2)u), I(W * |u|^2) \rangle dt \right| \lesssim \frac{\delta}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}}^5. \quad (32)$$

We have to show

$$D := \int_0^\delta \int_* M(\xi_1, \dots, \xi_5) \prod_{i=1}^5 \widehat{u}_i(\xi_i, t) d\xi_1 \dots d\xi_5 dt \lesssim \frac{\delta}{N^{2-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}[0, \delta]}$$

with

$$\begin{aligned}
M(\xi_1, \dots, \xi_5) &:= \frac{|m(\xi_1 + \xi_2 + \xi_3) - m(\xi_1)m(\xi_2)m(\xi_3)|}{m(\xi_1)m(\xi_2)m(\xi_3)\langle \xi_1 + \xi_2 \rangle^2} \\
&\quad \cdot \frac{m(\xi_4 + \xi_5)}{m(\xi_4)m(\xi_5)\langle \xi_4 + \xi_5 \rangle^2} \cdot \prod_{i=1}^5 |\xi_i|^{-1}.
\end{aligned}$$

We assume w.l.o.g. $N_1 \geq N_2$ and $N_4 \geq N_5$.

Case 1: $N \gg N_4 \geq N_5$.

a. $N_1 \geq N_2 \gtrsim N \gtrsim N_3$.

In this case we get

$$\begin{aligned}
D &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_2}{N}\right)^{\frac{1}{2}} \frac{1}{N_1 N_2} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^1 L_x^{\frac{6}{5}}} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \\
&\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^{\infty-}} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|D^{-1} u_3\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned} \tag{33}$$

b. $N_1, N_3 \gtrsim N \gg N_2$.

We have

$$\begin{aligned}
D &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_3}{N}\right)^{\frac{1}{2}} \frac{\delta}{N_1 N_3} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^2} \|u_3\|_{L_t^\infty L_x^2} \\
&\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

c. $N_1, N_2, N_3 \gtrsim N$.

This leads to

$$\begin{aligned}
D &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_2}{N}\right)^{\frac{1}{2}} \left(\frac{N_3}{N}\right)^{\frac{1}{2}} \frac{\delta}{N_1 N_2 N_3} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^2} \|u_3\|_{L_t^\infty L_x^2} \\
&\quad \cdot \|\langle D \rangle^{-2} (D^{-1} u_4 D^{-1} u_5)\|_{L_t^\infty L_x^\infty} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

Case 2: $N_4 \geq N_5 \gtrsim N$.

The additional factor $\left(\frac{N_4}{N}\right)^{\frac{1}{2}-} \left(\frac{N_5}{N}\right)^{\frac{1}{2}-}$ can be compensated by replacing the last factor in (33) by

$$N_4^{-\frac{1}{2}} N_5^{-\frac{1}{2}} \|D^{-\frac{1}{2}} u_4 D^{-\frac{1}{2}} u_5\|_{L_t^\infty L_x^{\frac{3}{2}}} \lesssim N_4^{-\frac{1}{2}} N_5^{-\frac{1}{2}} \|u_4\|_{L_t^\infty L_x^2} \|u_5\|_{L_t^\infty L_x^2}$$

leading to even an improved bound.

Case 3: $N_4 \geq N \geq N_5$.

We argue as before replacing the last factor in (33) by

$$N_4^{-\frac{1}{2}} \|D^{-\frac{1}{2}} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{2+}} \lesssim N_4^{-\frac{1}{2}} \|u_4\|_{L_t^\infty L_x^2} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}}.$$

This proves (32).

Next we claim

$$\left| \int_0^\delta \langle (W * |Iu|^2) Iu - I((W * |u|^2)u), I((W * Re u)u) \rangle dt \right| \lesssim \frac{\delta}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}}^5. \tag{34}$$

We again consider D with

$$\begin{aligned}
M(\xi_1, \dots, \xi_5) &:= \frac{|m(\xi_1 + \xi_2 + \xi_3) - m(\xi_1)m(\xi_2)m(\xi_3)|}{m(\xi_1)m(\xi_2)m(\xi_3)\langle \xi_1 + \xi_2 \rangle^2} \\
&\quad \cdot \frac{m(\xi_4 + \xi_5)}{m(\xi_4)m(\xi_5)\langle \xi_4 \rangle^2} \cdot \prod_{i=1}^5 |\xi_i|^{-1}.
\end{aligned}$$

We argue similarly as in the previous case and consider only the more difficult case $N_5 \geq N_4$.

Case 1: $N \gg N_5 \geq N_4$.

The last term in (33) with a suitable change of the Hölder exponent in the first factor can be replaced by

$$\begin{aligned} \|\langle D \rangle^{-2} D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^6} &\lesssim \|\langle D \rangle^{-2} D^{-1} u_4\|_{L_t^\infty L_x^\infty} \|D^{-1} u_5\|_{L_t^\infty L_x^6} \\ &\lesssim \|D^{-1} u_4\|_{L_t^\infty L_x^{6+}} \|u_5\|_{L_t^\infty L_x^2} \end{aligned}$$

leading to the same bound.

Case 2: $N_5 \geq N_4 \gtrsim N$.

The additional factor $(\frac{N_4}{N})^{\frac{1}{2}-} (\frac{N_5}{N})^{\frac{1}{2}-}$ is compensated by replacing the last factor in (33) by

$$\begin{aligned} N_4^{-\frac{1}{2}} N_5^{-\frac{1}{2}} \|\langle D \rangle^{-2} D^{-\frac{1}{2}} u_4 D^{-\frac{1}{2}} u_5\|_{L_t^\infty L_x^3} &\lesssim N_4^{-\frac{1}{2}} N_5^{-\frac{1}{2}} \|D^{-\frac{1}{2}} u_4\|_{L_t^\infty L_x^3} \|D^{-\frac{1}{2}} u_5\|_{L_t^\infty L_x^3} \\ &\lesssim N_4^{-\frac{1}{2}} N_5^{-\frac{1}{2}} \|u_4\|_{L_t^\infty L_x^2} \|u_5\|_{L_t^\infty L_x^2}. \end{aligned}$$

Case 3: $N_5 \geq N \geq N_4$.

Replace the last factor in (33) by

$$\begin{aligned} \|\langle D \rangle^{-2} D^{-1} u_4 D^{-1} u_5\|_{L_t^\infty L_x^{3+}} &\lesssim \|D^{-1} u_4\|_{L_t^\infty L_x^{6+}} \|D^{-1} u_5\|_{L_t^\infty L_x^6} \\ &\lesssim N_4^{0+} \|u_4\|_{L_t^\infty L_x^2} \|u_5\|_{L_t^\infty L_x^2}. \end{aligned}$$

Thus (34) is proven.

The next claim is

$$\left| \int_0^\delta \langle (W * |Iu|^2) - I(W * |u|^2), I((W * |u|^2)u) \rangle dt \right| \lesssim \frac{\delta}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}}^5. \quad (35)$$

We again consider D with

$$\begin{aligned} M(\xi_1, \dots, \xi_5) &:= \frac{|m(\xi_1 + \xi_2) - m(\xi_1)m(\xi_2)|}{m(\xi_1)m(\xi_2)\langle \xi_1 + \xi_2 \rangle^2} \\ &\quad \cdot \frac{m(\xi_3 + \xi_4 + \xi_5)}{m(\xi_3)m(\xi_4)m(\xi_5)\langle \xi_3 + \xi_4 \rangle^2} \cdot \prod_{i=1}^5 |\xi_i|^{-1}. \end{aligned}$$

We assume w.l.o.g. $N_2 \geq N_1$ and $N_3 \geq N_4$.

Case 1: $N \gg N_3, N_4, N_5$ ($\implies N_1 \sim N_2 \gtrsim N$).

In this case we get

$$\begin{aligned} D &\lesssim \left(\frac{N_1}{N}\right)^{\frac{1}{2}} \left(\frac{N_2}{N}\right)^{\frac{1}{2}} \frac{\delta}{N_1 N_2} \|\langle D \rangle^{-2} (u_1 u_2)\|_{L_t^\infty L_x^{\frac{6}{5}}} \|\langle D \rangle^{-2} (D^{-1} u_3 D^{-1} u_4)\|_{L_t^\infty L_x^{\infty-}} \\ &\quad \cdot \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|u_1 u_2\|_{L_t^\infty L_x^1} \|D^{-1} u_3 D^{-1} u_4\|_{L_t^\infty L_x^{3+}} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\ &\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}. \end{aligned}$$

Case 2: $N_3 \geq N \gg N_4, N_5$.

a. $N_3 \geq N \gg N_1$.

We get

$$\begin{aligned}
D &\lesssim \frac{\delta}{N_2 N_3} \|\langle D \rangle^{-2} (D^{-1} u_1 u_2)\|_{L_t^\infty L_x^2} \|\langle D \rangle^{-2} (u_3 D^{-1} u_4)\|_{L_t^\infty L_x^{3-}} \\
&\quad \cdot \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\
&\lesssim \frac{\delta}{N_2 N_3} \|D^{-1} u_1 u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3 D^{-1} u_4\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_2 \geq N_1 \gtrsim N$.

We get an additional factor $(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_2}{N})^{\frac{1}{2}-}$, which is acceptable, when one estimates as in a.

Case 3: $N_5 \geq N \gg N_3, N_4$ ($\implies N_2 \gtrsim N$).

a. $N_1 \leq N$.

We get

$$\begin{aligned}
D &\lesssim \frac{\delta}{N_2 N_5} \|\langle D \rangle^{-2} (D^{-1} u_1 u_2)\|_{L_t^\infty L_x^2} \|\langle D \rangle^{-2} (D^{-1} u_3 D^{-1} u_4)\|_{L_t^\infty L_x^\infty} \\
&\quad \cdot \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_2 N_5} \|D^{-1} u_1 u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|D^{-1} u_3 D^{-1} u_4\|_{L_t^\infty L_x^{3+}} \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_1 \geq N$.

We get an additional factor $(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_2}{N})^{\frac{1}{2}-}$, which can be compensated as in a.

Case 4: $N_3, N_4 \gtrsim N \gtrsim N_5$.

a. $N_2 \geq N \gg N_1$.

We get

$$\begin{aligned}
D &\lesssim (\frac{N_3}{N})^{\frac{1}{2}} (\frac{N_4}{N})^{\frac{1}{2}} \frac{\delta}{N_2 N_3 N_4} \|\langle D \rangle^{-2} (D^{-1} u_1 u_2)\|_{L_t^\infty L_x^2} \|\langle D \rangle^{-2} (u_3 u_4)\|_{L_t^\infty L_x^{3-}} \\
&\quad \cdot \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\
&\lesssim (\frac{N_3}{N})^{\frac{1}{2}} (\frac{N_4}{N})^{\frac{1}{2}} \frac{\delta}{N_2 N_3 N_4} \|D^{-1} u_1 u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_3 u_4\|_{L_t^\infty L_x^1} \|D^{-1} u_5\|_{L_t^\infty L_x^{6+}} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_1 \geq N$.

We get an additional factor $(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_2}{N})^{\frac{1}{2}-}$, which can be compensated by replacing the term $\|D^{-1} u_1 u_2\|_{L_t^\infty L_x^{\frac{3}{2}+}}$ in a. by

$$\|D^{-1} u_1 u_2\|_{L_t^\infty L_x^1} \lesssim \frac{1}{N_1} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2}.$$

Case 5: $N_3, N_5 \geq N \geq N_4$.

a. $N_2 \geq N \geq N_1$.

We have

$$\begin{aligned}
D &\lesssim \left(\frac{N_3}{N}\right)^{\frac{1}{2}} \left(\frac{N_5}{N}\right)^{\frac{1}{2}} \frac{\delta}{N_3 N_5} \|\langle D \rangle^{-2} (D^{-1} u_1 D^{-1} u_2)\|_{L_t^\infty L_x^{3+}} \|\langle D \rangle^{-2} (u_3 D^{-1} u_4)\|_{L_t^\infty L_x^{6-}} \\
&\quad \cdot \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{2-}} \|D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^{3+}} \|u_3 D^{-1} u_4\|_{L_t^\infty L_x^{\frac{3}{2}+}} \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{2-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_2 \geq N_1 \geq N$.

We get an additional factor $(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_2}{N})^{\frac{1}{2}-}$, which can be compensated by replacing the term $\|D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^{3+}}$ in a. by

$$\|D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^{1+}} \lesssim \frac{1}{N_1^{1-} N_2^{1-}} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2}.$$

Case 6: $N_3, N_4, N_5 \geq N$.

a. $N_2 \geq N \geq N_1$.

We have

$$\begin{aligned}
D &\lesssim \left(\frac{N_3}{N}\right)^{\frac{1}{2}} \left(\frac{N_4}{N}\right)^{\frac{1}{2}} \left(\frac{N_5}{N}\right)^{\frac{1}{2}} \frac{\delta}{N_3 N_4 N_5} \|\langle D \rangle^{-2} (D^{-1} u_1 D^{-1} u_2)\|_{L_t^\infty L_x^6} \\
&\quad \cdot \|\langle D \rangle^{-2} (u_3 u_4)\|_{L_t^\infty L_x^3} \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{3-}} \|D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^{3+}} \|u_3 u_4\|_{L_t^\infty L_x^1} \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{3-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

b. $N_2 \geq N_1 \geq N \implies N_{\min} \gtrsim N$.

We get an additional factor $(\frac{N_1}{N})^{\frac{1}{2}-} (\frac{N_2}{N})^{\frac{1}{2}-} \lesssim (\frac{N}{N})^{1-}$ leading to

$$\begin{aligned}
D &\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|\langle D \rangle^{-2} (u_1 D^{-1} u_2)\|_{L_t^\infty L_x^6} \|\langle D \rangle^{-2} (u_3 u_4)\|_{L_t^\infty L_x^3} \\
&\quad \cdot \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta}{N_{\max}^{0+} N^{4-}} \|u_1 D^{-1} u_2\|_{L_t^\infty L_x^{\frac{3}{2}}} \|u_3 u_4\|_{L_t^\infty L_x^1} \|u_5\|_{L_t^\infty L_x^2} \\
&\lesssim \frac{\delta(1 \wedge N_{\min}^{0+})}{N_{\max}^{0+} N^{4-}} \prod_{i=1}^5 \|u_i\|_{X^{0, \frac{1}{2}+}}.
\end{aligned}$$

which completes the proof of (35).

Next we want to prove the following estimate:

$$|\int_0^\delta \langle (W * Re Iu) Iu - I((W * Re u)u), I((W * |u|^2)u) \rangle dt| \lesssim \frac{\delta}{N^{2-}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}}^5. \quad (36)$$

We again consider D with

$$\begin{aligned}
M(\xi_1, \dots, \xi_5) &:= \frac{|m(\xi_1 + \xi_2) - m(\xi_1)m(\xi_2)|}{m(\xi_1)m(\xi_2)\langle \xi_1 \rangle^2} \\
&\quad \cdot \frac{m(\xi_3 + \xi_4 + \xi_5)}{m(\xi_3)m(\xi_4)m(\xi_5)\langle \xi_3 + \xi_4 \rangle^2} \cdot \prod_{i=1}^5 |\xi_i|^{-1},
\end{aligned}$$

and treat only the more difficult case $N_2 \geq N_1$, and assume w.l.o.g. $N_3 \geq N_4$. We consider the same cases as for (35).

In case 1 we replace $\|\langle D \rangle^{-2}(u_1 u_2)\|_{L_t^\infty L_x^{\frac{6}{5}}}$ by

$$\|\langle D \rangle^{-2} u_1 u_2\|_{L_t^\infty L_x^{\frac{6}{5}}} \lesssim \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2}.$$

In the cases 2,3 and 4a we replace $\|\langle D \rangle^{-2}(D^{-1} u_1 u_2)\|_{L_t^\infty L_x^2}$ by

$$\|\langle D \rangle^{-2} D^{-1} u_1 u_2\|_{L_t^\infty L_x^2} \lesssim \|D^{-1} u_1\|_{L_t^\infty L_x^{6+}} \|u_2\|_{L_t^\infty L_x^2}.$$

In case 4b. estimate

$$\|\langle D \rangle^{-2} D^{-1} u_1 u_2\|_{L_t^\infty L_x^2} \lesssim \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2}.$$

Case 5a is essentially unchanged, whereas in Case 5b replace

$\|\langle D \rangle^{-2}(D^{-1} u_1 D^{-1} u_2)\|_{L_t^\infty L_x^{3+}} \|\langle D \rangle^{-2}(u_3 D^{-1} u_4)\|_{L_t^\infty L_x^{6-}}$ by

$$\begin{aligned} & \|\langle D \rangle^{-2} D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^{2+}} \|\langle D \rangle^{-2}(u_3 D^{-1} u_4)\|_{L_t^\infty L_x^{\infty-}} \\ & \lesssim \|D^{-1} u_1\|_{L_t^\infty L_x^2} \|D^{-1} u_2\|_{L_t^\infty L_x^{2+}} \|u_3 D^{-1} u_4\|_{L_t^\infty L_x^{\frac{3}{2}+}} \\ & \lesssim \frac{1}{N_1 N_2^{1-}} \|u_1\|_{L_t^\infty L_x^2} \|u_2\|_{L_t^\infty L_x^2} \|u_3\|_{L_t^\infty L_x^2} \|D^{-1} u_4\|_{L_t^\infty L_x^{6+}}. \end{aligned}$$

In case 6a estimate

$$\begin{aligned} \|\langle D \rangle^{-2} D^{-1} u_1 D^{-1} u_2\|_{L_t^\infty L_x^6} & \lesssim \|\langle D \rangle^{-2} D^{-1} u_1\|_{L_t^\infty L_x^\infty} \|D^{-1} u_2\|_{L_t^\infty L_x^6} \\ & \lesssim \|D^{-1} u_1\|_{L_t^\infty L_x^{6+}} \|u_2\|_{L_t^\infty L_x^2}. \end{aligned}$$

Similarly case 6b can be handled, so that (36) is complete.

The forth and third order terms in $|\langle F(Iu) - IF(u), IF(u) \rangle|$ turn out to be less critical. We omit any detailed calculations here and just refer to the recent paper of the author [Pe], where the following estimates were given even under the weaker assumption $|\widehat{W}(\xi)| \lesssim 1$ (compared to the property $|\widehat{W}(\xi)| \lesssim \langle \xi \rangle^{-2}$ which we have in the present study). We have to remark that the assumption $s \geq \frac{3}{4}$ in that paper is not really necessary for these forth and third order terms, but could be replaced by $s > \frac{1}{2}$, because the factors $(\frac{N_i}{N})^{\frac{1}{4}}$ can everywhere be replaced by $(\frac{N_i}{N})^{\frac{1}{2}-}$ without significance for the results. We have ([Pe], section 4.6, 4.7 and 4.8):

$$\begin{aligned} \int_0^\delta |I(u^3) - (Iu)^3, Iu| dt & \lesssim N^{-3+} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^4, \\ \int_0^\delta |I(u^2) - (Iu)^2, (Iu)^2| dt & \lesssim N^{-3+} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^4, \\ \int_0^\delta |I(u^2) - (Iu)^2, Iu| dt & \lesssim N^{-\frac{5}{2}+} \delta^{\frac{1}{2}} \|\nabla Iu\|_{X^{0, \frac{1}{2}+}[0, \delta]}^3. \end{aligned}$$

Summarizing all our estimates in this chapter we arrive at (17).

REFERENCES

- [ABJ] A. Aftalion, X. Blanc and R.L. Jerrard: *Nonclassical rotational inertia of a supersolid*. Phys. Review Letters 99 (2007), 135301-4
- [BS] F. Bethuel and J.C. Saut: *Travelling waves for the Gross-Pitaevskii equation I*. Ann. I. H. Poincaré Phys. Théor. 70 (1999), 147-238
- [CH] T. Cazenave and A. Haraux: *An introduction to semilinear evolution equations*. Oxford science publications 1998
- [CKSTT] J. Colliander, M. Keel, G. Staffilani, H. Takaoka and T. Tao: *Almost conservation laws and global rough solutions to a nonlinear Schrödinger equation*. Math. Res. Letters 9 (2002), 659-682

- [C] C. Coste: *Nonlinear Schrödinger equation and superfluid hydrodynamics*. Eur. Phys. J. B Condens. Matter Phys. (1998), 245-253
- [Ga] C. Gallo: *The Cauchy problem for defocusing nonlinear Schrödinger equations with non-vanishing initial data at infinity*. Comm. Part. Diff. Equa. 33 (2008), 729-771
- [Ge] P. Gérard: *The Cauchy problem for the Gross-Pitaevskii equation*. Ann. I. H. Poincaré Anal. Non-linéaire 23 (2006), 765-779
- [GTV] J. Ginibre, Y. Tsutsumi and G. Velo: *On the Cauchy problem for the Zakharov system*. J. Funct. Analysis 151 (1997), 384-436
- [Gr] E.P. Gross: *Hydrodynamics of a Superfluid Condensate*. J. Math. Phys. 4 (1963), 195-207
- [G] A. Grünrock: *New applications of the Fourier restriction norm method to wellposedness problems for nonlinear evolution equations*. Dissertation Univ. Wuppertal 2002, <http://elpub.bib.uni-wuppertal.de/servlets/DocumentServlet?id=254>
- [JPR] Ch. Josserand, Y. Pomeau and S. Rica: *Coexistence of ordinary elasticity and superfluidity in a model of a defect-free supersolid*. Phys. Review Letters 98 (2007), 195301-4
- [JPRo] C.A. Jones, S.J. Putterman, P.H. Roberts: *Motion in a Bose condensate V. Stability of solitary wave solutions of non-linear Schrödinger equations in two and three dimensions*. J. Phys. A, Math. Gen. 19 (1986), 2991-3011
- [JR] C.A. Jones and P.H. Roberts: *Motions in a Bose condensate IV. Axisymmetric solitary waves*. J. Phys. A, Math. Gen. 15 (1982), 2599-2619
- [KT] M. Keel and T. Tao: *Endpoint Strichartz estimates*. Amer. J. Math. 120 (1998), 955-98
- [KL] Y.S. Kivshar and B. Luther-Davies: *Dark optical solitons: physics and applications*. Phys. Rep. 298 (1998), 81-197
- [L] A. de Laire: *Global well-posedness for a nonlocal Gross-Pitaevskii equation with non-zero condition at infinity*. Comm in PDE 35 (2010), 2021-2058
- [Pe] H. Pecher: *Unconditional global well-posedness for the 3D Gross-Pitaevskii equation for data without finite energy*. arXiv:1201.3777
- [P] L.P. Pitaevskii: *Vortex lines in an imperfect Bose gas*. Soviet Physics JETP 13 (1961), 451-454
- [PR] Y. Pomeau and S. Rica: *Model of superflow with rotons*. Phys. Review Letters 71 (1993), 247-250
- [SK] V.S. Shchesnovich and R.A. Kraenkel: *Vortices in lonlocal Gross-Pitaevskii equation*. J. Phys. A 37 (2004), 6633-6651

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